

S.Bernuzzi CINECA Nov 15th 2017

Numerical relativity in the gravitational-wave astronomy era



First numerical relativity simulation of neutron star merger with precessing spins: the double pulsar case



Viz by T.Dietrich

Baryon mass density

GWs: Tiny signatures of extreme events

Collision of neutron stars [Mass~1.4 Msun, Radius~10 km]:



First numerical relativity simulation of neutron star merger with precessing spins: the double pulsar case



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Weyl curvature scalar

What can we learn from neutron star mergers?

FUNDAMENTAL PHYSICS Strong-field tests GR (dynamics) Structure of bulk matter at supranuclear densities Heavy elements nucleosynthesis







ASTROPHYSICS (Multi-messenger) Origin of gamma-ray burst Origin of kilonovae, site for r-processes





COSMOGRAPHY Measure Hubble constant *Standard sirens*, Calibrate cosmic distance ladder

Fundamental physics

Constraining the Equation of State of matter at supranuclear densities



Different EOS → different star's structure





Binary neutron star mergers

Example: observing tidal effects in GWs tells us about the neutron star matter



Tides determine the wave's phase during merger

Example: observing tidal effects in GWs tells us about the neutron star matter



Example: observing tidal effects in GWs tells us about the neutron star matter



The GW spectrum of binary neutron stars



- Faithful and **complete waveform model** (*inspiral+merger+postmerger*)
- Coverage of the **parameter space** (mass, spins, EOS, ...)
- Precise prediction of the merger remnant (e.g. collapse, black hole)

Methods for the GR 2-body problem



$$\begin{split} \partial_t \tilde{\Gamma}^i &= -2\,\tilde{A}^{ij}\,\partial_j \alpha + 2\,\alpha \left[\tilde{\Gamma}^i{}_{jk}\,\tilde{A}^{jk} - \frac{3}{2}\,\tilde{A}^{ij}\,\partial_j \ln(\chi) \right. \\ &\left. -\frac{1}{3}\,\tilde{\gamma}^{ij}\,\partial_j(2\,\hat{K} + \Theta) - 8\,\pi\,\tilde{\gamma}^{ij}\,S_j \right] + \tilde{\gamma}^{jk}\,\partial_j\partial_k\beta \\ &\left. + \frac{1}{3}\,\tilde{\gamma}^{ij}\partial_j\partial_k\beta^k + \beta^j\,\partial_j\tilde{\Gamma}^i - (\tilde{\Gamma}_d)^j\,\partial_j\beta^i \right. \\ &\left. + \frac{2}{3}\,(\tilde{\Gamma}_d)^i\,\partial_j\beta^j - 2\,\alpha\,\kappa_1\,\left[\tilde{\Gamma}^i - (\tilde{\Gamma}_d)^i\right], \right. \\ &\left. \partial_t\Theta = \frac{1}{2}\,\alpha\left[R - \tilde{A}_{ij}\,\tilde{A}^{ij} + \frac{2}{3}\,(\hat{K} + 2\,\Theta)^2\right] \\ &\left. - \alpha\left[8\,\pi\,\rho + \kappa_1\,(2 + \kappa_2)\,\Theta\right] + \beta^i\partial_i\Theta\,, \end{split}$$

GR Formulation and Cauchy problem + GR hydrodynamics



Coordinates and Singularities

Numerical relativity in a nutshell

Numerical methods for PDEs on adaptive grids



High-performance-computing (HPC)



Numerical relativity: Cauchy problem in GR





- 3+1 formulation (hyperboloidal slices?)
- Initial data (Lichnerowic, York, ...)
- Evolution schemes (GHG; ADM \rightarrow BSSN, Z4c)
- Well posedness (Choquet-Bruhat; Friedrich; Gundlach&Martin-Garcia) [*need gauge fix*]

$$\begin{split} \partial_t \chi &= \frac{2}{3} \chi [\alpha(\hat{K} + 2\Theta) - D_i \beta^i], \\ \partial_t \tilde{\gamma}_{ij} &= -2\alpha \tilde{A}_{ij} + \beta^k \tilde{\gamma}_{ij,k} + 2\tilde{\gamma}_{k(i} \beta^k_{,j)} - \frac{2}{3} \tilde{\gamma}_{ij} \beta^k_{,k} , \\ \partial_t \hat{K} &= -D^i D_i \alpha + \alpha [\tilde{A}_{ij} \tilde{A}^{ij} + \frac{1}{3} (\hat{K} + 2\Theta)^2] \\ &+ 4\pi \alpha [S + \rho_{\rm ADM}] + \alpha \kappa_1 (1 - \kappa_2) \Theta + \beta^i \hat{K}_{,i} \\ \partial_t \tilde{A}_{ij} &= \chi [-D_i D_j \alpha + \alpha (R_{ij} - 8\pi S_{ij})]^{\rm tf} \\ &+ \alpha [(\hat{K} + 2\Theta) \tilde{A}_{ij} - 2\tilde{A}^k_i \tilde{A}_{kj}] \\ &+ \beta^k \tilde{A}_{ij,k} + 2\tilde{A}_{k(i} \beta^k_{,j)} - \frac{2}{3} \tilde{A}_{ij} \beta^k_{,k} \\ \partial_t \tilde{\Gamma}^i &= -2\tilde{A}^{ij} \alpha_{,j} + 2\alpha [\tilde{\Gamma}^i_{jk} \tilde{A}^{jk} - \frac{3}{2} \tilde{A}^{ij} \ln(\chi)_{,j} \\ &- \frac{1}{3} \tilde{\gamma}^{ij} (2\hat{K} + \Theta)_{,j} - 8\pi \tilde{\gamma}^{ij} S_j] + \tilde{\gamma}^{jk} \beta^i_{,jk} \\ &+ \frac{1}{3} \tilde{\gamma}^{ij} \beta^k_{,kj} + \beta^j \tilde{\Gamma}^i_{,j} - \tilde{\Gamma}_d{}^j \beta^i_{,j} + \frac{2}{3} \tilde{\Gamma}_d{}^i \beta^j_{,j} \\ &- 2\alpha \kappa_1 (\tilde{\Gamma}^i - \tilde{\Gamma}_d{}^i), \\ \partial_t \Theta &= \alpha [\frac{1}{2} R - \frac{1}{2} \tilde{A}_{ij} \tilde{A}^{ij} + \frac{1}{3} (\hat{K} + 2\Theta)^2 \\ &- 8\pi \rho_{\rm ADM} - \kappa_1 (2 + \kappa_2) \Theta] + \mathcal{L}_\beta \Theta. \end{split}$$

[SB, Hilditch arXiv:0912.2920]

Numerical relativity: singularities & crash tests





- Crash at tau=pi (geodesic slicing)
- Lapse collapse, slice stretching (1+log, shift=0)
 e.g. [Bruegmann arXiv:9912009]

Numerical relativity: singularities & coordinates



[Thierfelder, SB, Hilditch, Bruegmann, Rezzolla arXiv:1012.3703]

Numerical relativity: numerical methods (some)

Adaptive mesh refinement (AMR) \rightarrow resolve multiple scales





Grid based, AMR Berger-Oliger Method of line w\ Runge-Kutta (Subcycling) Finite differencing and finite volumes Numerical relativity specs

- R.H.S. complexity (derivatives and contractions)
- \rightarrow stencil ("horizontal") + pointwise ("vertical") ops
- High-order operators (large 3D stencils > 5 pts/direction)
- \rightarrow communication overhead for distributed computations
- Memory: >~ O(100) 3D grid function per time level

Improved NR GW with high-order WENO schemes

[SB,Dietrich PRD94 064062 (2016)]



- Robust convergence assessment (although not 5th order)
- Large resolution span (64³-192³), no alignment
- Error budget: significant improvement wrt FV schemes

See also [SB+ arXiv:1205.3403] [Radice+ arxiv:1306.6052]

First waveform model for inspiral → merger

[SB,Nagar,Dietrich,Damour PRL 114 (2015)]



- Effective-one-body model with tides, GSF Resummed approach [Bini+ 2014]
- Valid from low frequencies to merger, PREDICT the merger waveform
- Accuracy: uncertainties of the numerical data (improve simulations!)

See [Hinderer+ PRL 116 (2016)] for an alternative approach

Spins & tides during merger: phasing

[Dietrich, SB, Ujevic, Tichy PRD 95 (2017)]



Frequency-domain tidal wf approximant

[Dietrich, SB, Tichy arXiv:1706.02969]

First NR-based tidal approximant Fast, flexible, accurate



Exploring the BNS parameter space

Largest exploration of parameter space in strong-field regime available to date



SB&Dietrich PRD94 (2016), Dietrich+ PRD95 024029 (2017), Radice+ ApJL 842 (2017),]

HPC time usage

Usage in 2015/2016



Production codes & parallelization

- BAM [Bruegmann (Jena) + Tichy (Florida Atlantic) + Bernuzzi (Parma) and others]
- THC [Radice (Princeton), based on Cactus & ET/CTGamma]

- MPI: domain decomposition for each level
- OMP threads on MPI job
- None or poor vectorization
- Numerical relativity specs:
 - − R.H.S. complexity (derivatives and contractions) → stencil ("horizontal") + pointwise ("vertical") ops
 - High-order operators (large 3D stencils > 5 pts/direction) \rightarrow communication overhead for distributed computations
 - Memory: >~ O(100) 3D grid function per time level

No NR production code scales to >~ 10k cores

Porting to KNL: first strong scaling results



BAM (MPI+OMP) Numerical relativity, Compact binaries simulations FISH-ASL (MPI) Newtonian gravity + hydro Supernova core-collapse and disk winds

NOTE: no attempt to optimize code, just re-compile and run

Linearized Einstein

The spacetime metric reads

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu},$$

where $g_{\mu\nu}$ is the background metric and $h_{\mu\nu}$ is a small perturbation. We adopt the following gauge

$$g_{00} = 1, \quad g_{0i} = 0,$$

so that the linearized Einstein equations read

$$\partial_t^2 h_{ij} = \eta^{kl} \partial_k \partial_l h_{ij}.$$

The equations are solved in first order in time form

$$egin{aligned} \partial_t h_{ij} &= -2K_{ij}\ R_{ij} &= -rac{1}{2}\eta^{kl}\partial_k\partial_l h_{ij},\ \partial_t K_{ij} &= R_{ij}. \end{aligned}$$
 Tensor we

Tensor wave equations

Operations

 $R_{ij} = -\frac{1}{2} \eta^{kl} \partial_k \partial_l h_{ij}$

Tensor contractions

R23

-0.5*deldelg1123 + gamma113*gammado211 + gamma213*gammado212 +
(gamma212 + gamma313)*gammado213 + gamma112*(gammado113 + gammado311) + gamma212*gammado312 +
2.*gamma312*gammado313)*ginv11 +
(-deldelg1223 + gamma123*gammado211 + (gamma113 + gamma223)*gammado212 +
(gamma222 + gamma323)*gammado213 + gamma213*gammado222 +
(gamma212 + gamma313)*gammado223 + gamma122*(gammado112 + gammado211) + gamma2222*gammado212 +
gammall2*(gammadol23 + gammadol22) + gamma212*gammadol22 +
2.*(gamma322*gammado313 + gamma312*gammado323))*ginv12 +
(-deldelg1323 + gamma133*gammado211 + gamma233*gammado212 +
(gamma113 + gamma223 + gamma333)*gammado213 + gamma213*gammado223 +
(gamma212 + gamma313)*gammado233 +
gammal23*(gammadoll3 + gammado311) + gamma223*gammado312 +
gdmmd112*(gdmmd00133 + gdmmd00313) + gdmmd212*gdmmd00323 + 2 */aamma323*aammado313 + aamma312*aammado333))*ainv13 +
(-0.5*deldelg2223 + gamma123*gammado212 + gamma223*gammado222 +
(gamma222 + gamma323)*gammado223 +
gamma122*(gammado123 + gammado312) + gamma222*gammado322 +
2.*gamma322*gammado323)*ginv22 +
(-deldelg2323 + gamma133*gammado212 + gamma233*gammado222 +
(2.39) (animazzo + gammazzo) + gammado2zo + (animazzo + animazzo + animazo + animazzo + animazzo
gammal23*(gammadol23 + gammado213 + gammado312) +
gamma122*(gammado133 + gammado313) + gamma223*gammado322 +
(gamma222 + 2.*gamma323)*gammado323 + 2.*gamma322*gammado333)*ginv23 +
(-0.5*deldelg3323 + gamma133*gammado213 + gamma233*gammado223 +
(gamma223 + gamma533)*gammad0233 + gamma123*(gammad0133 + gammad0313) + gamma223*gammad0323 +
2.*qamma323*qammado333)*qinv33 +
0.5*((gammado213 + gammado312)*Gfromg1 +
(gammado223 + gammado322)*Gfromg2 + (gammado233 + gammado323)*Gfromg3 +
delG31*g12[ijk] + delG21*g13[ijk] + delG32*g22[ijk] +
$(aeio_{22} + aeio_{33}) + g_{23}[1]K] + aeio_{23} + g_{33}[1]K])$

Partial derivatives (3D) *



U[i,j,k] = C[0]*V[i,j,k]; for (r=1; r<=4; r++) U[i,j,k]+=C[r]*(V[i-r,j,k]+V[i,j-r,k]+V[i,j,k-r]+ V[i+r,j,k]+V[i,j+r,k]+V[i,j,k+r]);

pernuzzi@ThinkPad-T440s:~/Codes/BAM/src/projects/z4\$ wc -l z4_rhs_movpunc_N.c 10854 z4 rhs movpunc N.c

*[https://software.intel.com/en-us/articles/3d-finite-differences-on-multi-core-processors/]

An anisotropic and three-components kilonova counterpart of GW170817



Joint constraint on the neutron star equation of state from multimessenger observations



Summary

- Numerical relativity is key for
 - Waveform modeling
 - Exploting multi-messenger astronomy info
- No standard but dedicated codes
- New challenges need
 - Improve performances
 - New solutions