

**BLOW-UP OF COMPLEX SOLUTIONS
OF THE 3-d NAVIER-STOKES EQUATIONS
AND BEHAVIOR OF RELATED REAL SOLUTIONS**

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1. INTRODUCTION

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$$\frac{\partial \mathbf{u}}{\partial t} + \sum_{j=1}^3 u_j \frac{\partial}{\partial x_j} \mathbf{u} = \Delta \mathbf{u} - \nabla p, \quad \mathbf{x} = (x_1, x_2, x_3) \in \mathbb{R}^3.$$

$$\nabla \cdot \mathbf{u} = 0, \quad \mathbf{u}(\cdot, 0) = \mathbf{u}_0.$$

$\mathbf{u} : \mathbb{R}^3 \times [0, \infty) \rightarrow \mathbb{R}^3$ is the velocity field, p is the pressure and we assume for the viscosity $\nu = 1$ (always possible by rescaling).

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In spite of considerable progress, it is still unknown whether there are initial conditions for which the solution becomes singular in a finite time (global regularity problem).

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where a blowup is proved for a modified NS equations, which preserve the *energy identity*

$$E(t) + \int_0^t S(\tau) d\tau = E(0),$$

$E(t)$ is the total energy and $S(t)$ the total enstrophy

$$E(t) = \frac{1}{2} \int_{\mathbb{R}^3} |\mathbf{u}(\mathbf{x}, t)|^2 d\mathbf{x}, \quad S(t) = \int_{\mathbb{R}^3} |\nabla \times \mathbf{u}(\mathbf{x}, t)|^2 d\mathbf{x}.$$

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Poisson equation with Dirichlet boundary condition.

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Using the Fourier inversion theorem:

$$\mathbf{u}(\mathbf{x}, t) = -i \int_{\mathbb{R}^3} e^{-i\langle \mathbf{k}, \mathbf{x} \rangle} \mathbf{v}(\mathbf{k}, t) d\mathbf{k}$$

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$$-\nabla \cdot p(\mathbf{x}, t) = i \int_{\mathbb{R}^3} e^{-i\langle \mathbf{k}, \mathbf{x} \rangle} \left(\int_{\mathbb{R}^3} \langle \mathbf{v}(\mathbf{k} - \mathbf{k}', t), \mathbf{k}' \rangle \frac{(\mathbf{k} \cdot \mathbf{v}(\mathbf{k}', t))}{|\mathbf{k}|^2} \mathbf{k} d\mathbf{k}' \right) d\mathbf{k}$$

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Integrating:

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$$\mathbf{v}(\mathbf{k}, t) = e^{-tk^2} \mathbf{v}_0(\mathbf{k}) + \int_0^t e^{-(t-s)k^2} \left(\int_{\mathbb{R}^3} \langle \mathbf{v}(\mathbf{k} - \mathbf{k}', s), \mathbf{k} \rangle P_{\mathbf{k}} \mathbf{v}(\mathbf{k}', s) d\mathbf{k}' \right) ds,$$

using the linearity of the integral:

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The antitransform $\mathbf{u}(\mathbf{x}, t)$ is complex in general.

Integrating:

$$\mathbf{v}(\mathbf{k}, t) = e^{-t\mathbf{k}^2} \mathbf{v}_0(\mathbf{k}) + \int_0^t e^{-(t-s)\mathbf{k}^2} \left(\int_{\mathbb{R}^3} \langle \mathbf{v}(\mathbf{k} - \mathbf{k}', s), \mathbf{k} \rangle P_{\mathbf{k}} \mathbf{v}(\mathbf{k}', s) d\mathbf{k}' \right) ds,$$

using the linearity of the integral:

$$\int_{\mathbb{R}^3} \langle \mathbf{v}(\mathbf{k} - \mathbf{k}', t), \mathbf{k} \rangle P_{\mathbf{k}} \mathbf{v}(\mathbf{k}', t) d\mathbf{k}' = P_{\mathbf{k}} \int_{\mathbb{R}^3} \langle \mathbf{v}(\mathbf{k} - \mathbf{k}', t), \mathbf{k} \rangle \mathbf{v}(\mathbf{k}', t) d\mathbf{k}'$$

then

$$\mathbf{v}(\mathbf{k}, t) = e^{-t\mathbf{k}^2} \mathbf{v}_0(\mathbf{k}) + \int_0^t e^{-(t-s)\mathbf{k}^2} \left(P_{\mathbf{k}} \int_{\mathbb{R}^3} \langle \mathbf{v}(\mathbf{k} - \mathbf{k}', s), \mathbf{k} \rangle \mathbf{v}(\mathbf{k}', s) d\mathbf{k}' \right) ds,$$

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If however $\mathbf{v}(\mathbf{k}, t)$ is odd in \mathbf{k} then $\mathbf{u}(\mathbf{x}, t)$ is real and odd in \mathbf{x} .

$$\int_{\mathbb{R}^3} \langle \mathbf{v}(\mathbf{k}-\mathbf{k}', s), \mathbf{k} \rangle v_j(\mathbf{k}', s) d\mathbf{k}' = \sum_{i=1}^3 k_i \int_{\mathbb{R}^3} v_i(\mathbf{k}-\mathbf{k}', s) v_j(\mathbf{k}', s) d\mathbf{k}' =$$

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the evolution equation became

$$\mathbf{v}(\mathbf{k}, t+\Delta t) = e^{-(t+\Delta t)\mathbf{k}^2} \mathbf{v}_0(\mathbf{k}) + \int_0^{t+\Delta t} e^{-(t+\Delta t-s)\mathbf{k}^2} \mathbf{f}(\mathbf{v}, \mathbf{k}, s) ds =$$

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where $\beta_I = 1$, $\beta_{II} = \frac{1}{2}$ and $C_E^{(\alpha)}$, $C_S^{(\alpha)}$ are constants.

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- iii) The solutions converge point-wise in \mathbf{k} -space as $t \uparrow \tau$, while $E(t)$ and $S(t)$ diverge as inverse powers of $\tau - t$.

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Computer simulations for the Li-Sinai solutions were first performed by Arnol'd and Khokhlov in 2009. However, due to computational limitations, they could only get a qualitative description of the blow-up.

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We get solutions of type *I* for the initial data \mathbf{v}_0^+ , and of type *II* (alternating signs) for \mathbf{v}_0^- .

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The mesh R in \mathbf{k} -space is taken with step 1:

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For the description of the behavior we use the energy and enstrophy marginals in \mathbf{k} along the main axis

$$E_3(k_3, t) = \frac{1}{2} \int_{\mathbb{R} \times \mathbb{R}} dk_1 dk_2 |\mathbf{v}(\mathbf{k}, t)|^2,$$

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The marginals in \mathbf{x} -space are denoted $\tilde{E}_j(x_j, t)$, $\tilde{S}_j(x_j, t)$, $j = 1, 2, 3$, with

$$\tilde{E}_3(k_3, t) = \frac{1}{2} \int_{\mathbb{R} \times \mathbb{R}} dx_1 dx_2 |\mathbf{u}(\mathbf{x}, t)|^2,$$

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- The total enstrophy $S(t)$ starts growing much earlier than the total energy $E(t)$.
- The solution of type I blow up much earlier than the solutions of type II with the same initial energy and same a .

3. Li-Sinai solutions: simulations. The fixed point $\mathbf{H}^{(0)}$.

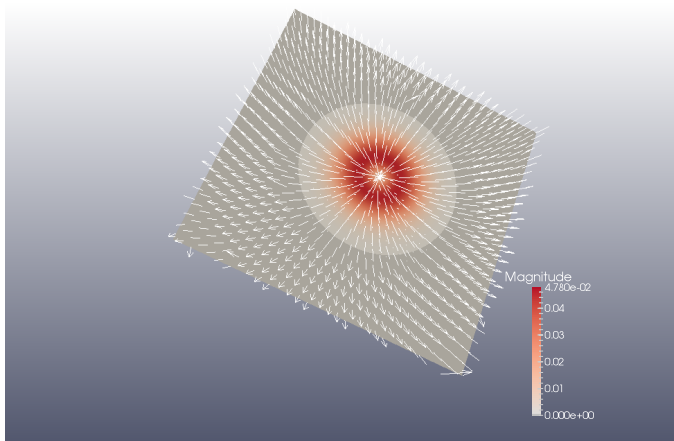


Figure 1: Type II, $a = 20$, $E_0 = 5 \times 10^4$. The arrows indicate the direction of $\mathbf{v}(\mathbf{k}, t)$ on a regular point lattice on a section of the plane $k_3 = 100$ with sides of length 100, $t = 1521 \times 10^{-7}$. Simulation range $k_3 \in [-19, 2528]$. Magnitude refers to $|\mathbf{v}(\mathbf{k}, t)|$. In the grey external region $|\mathbf{v}(\mathbf{k}, t)| < 10^{-6}$.

3. Li-Sinai solutions: simulations. Oscillations type I.

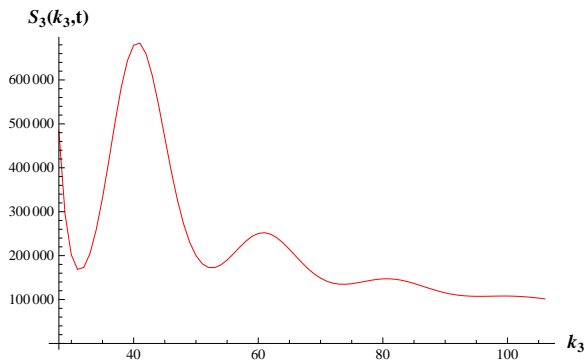


Figure 2: Type I, $a = 20$, $E_0 = 5 \times 10^4$. Enstrophy marginal density $S_3(k_3, t)$ at the beginning of the blow-up. $t = 900 \times 10^{-7}$. Simulation range $k_3 \in [-19, 2528]$.

3. Li-Sinai solutions: simulations. Oscillations type II.

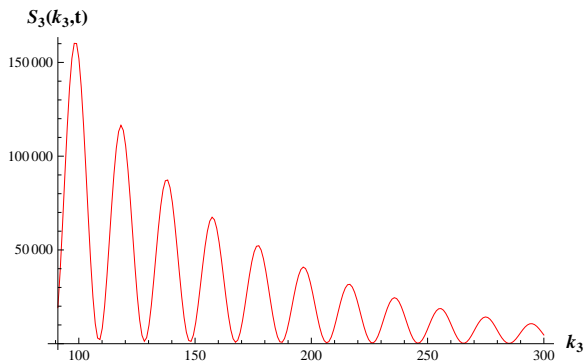


Figure 3: Type II, $a = 20$, $E_0 = 5 \times 10^4$. Entrophy marginal density $S_3(k_3, t)$ at the beginning of the blow-up. $t = 1125 \times 10^{-7}$. The zeroes are approximately periodic with period a . Simulation range $k_3 \in [-19, 2528]$.

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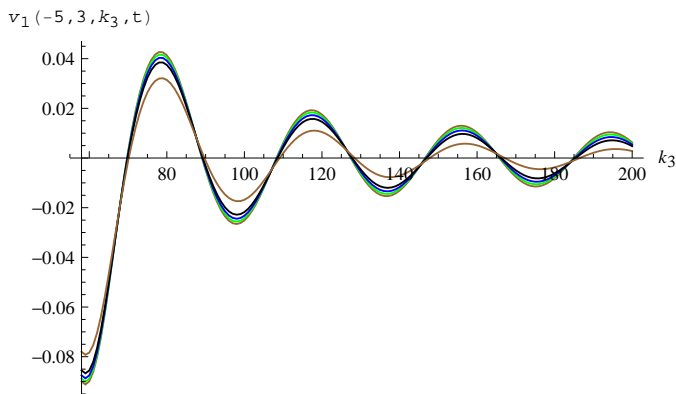


Figure 4: Type II, $a = 20$, $E_0 = 5 \times 10^4$. $\mathbf{v}_1(\mathbf{k}, t)$ vs k_3 for k_1, k_2 fixed, at the times $t \times 10^7 = 1342, 1500, 1544, 1574, 1600$. The amplitudes increase as time grows, and tend to a limit. Simulation range $k_3 \in [-19, 3028]$.

3. Li-Sinai solutions: simulations. Type I: compared growth.

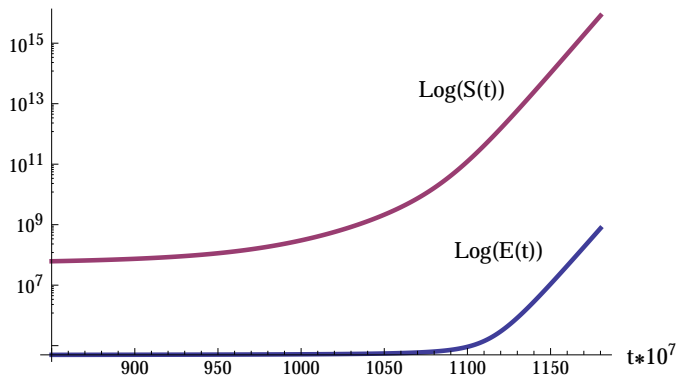


Figure 5: Type I, $a = 20$, $E_0 = 5 \times 10^4$. Compared growth of the total enstrophy $S(t)$ and the total energy $E(t)$. Simulation range $k_3 \in [-19, 2528]$.

3. Li-Sinai solutions: simulations. Type II: compared growth

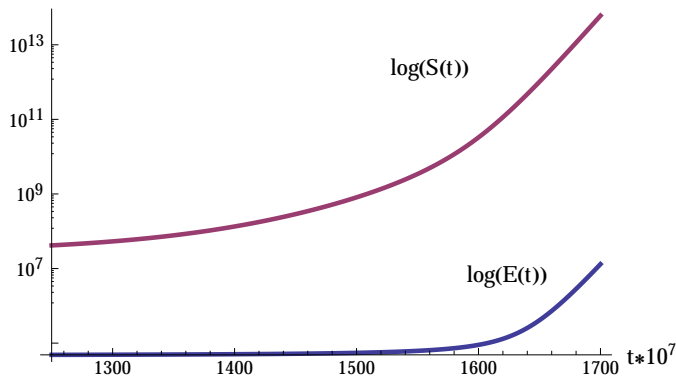


Figure 6: Type II, $a = 20$, $E_0 = 5 \times 10^4$. Compared growth of the total entropy $S(t)$ and the total energy $E(t)$. Simulation range $k_3 \in [-19, 2528]$.

3. Li-Sinai solutions: simulations. Enstrophy distribution

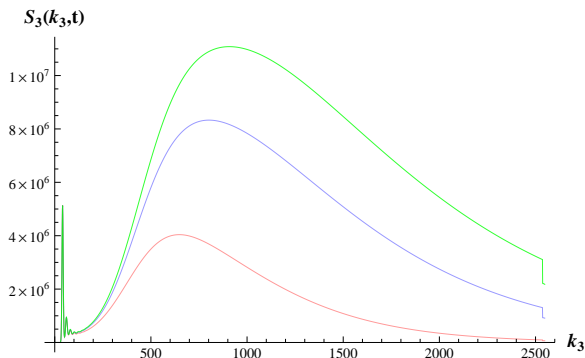


Figure 7: Type I, $a = 20$, $E_0 = 5 \times 10^4$. Plot of the marginal enstrophy density $S_3(k_3, t)$ on the whole simulation range $-19 \leq k_3 \leq 2528$, at $t \cdot 10^7 = 1060, 1075, 1080$.

3. Li-Sinai solutions: simulations. Enstrophy distribution

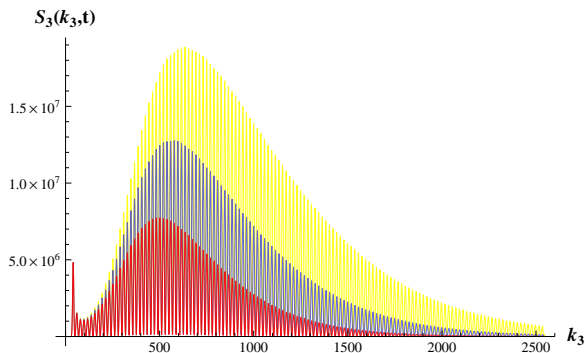


Figure 8: *Type II*, $a = 20$, $E_0 = 5 \times 10^4$. Plot of the marginal enstrophy density $S_3(k_3, t)$ on the whole simulation range $-19 \leq k_3 \leq 2528$, at $t \cdot 10^7 = 1521, 1544, 1560$.

3. Li-Sinai solutions: simulations. Enstrophy distribution.

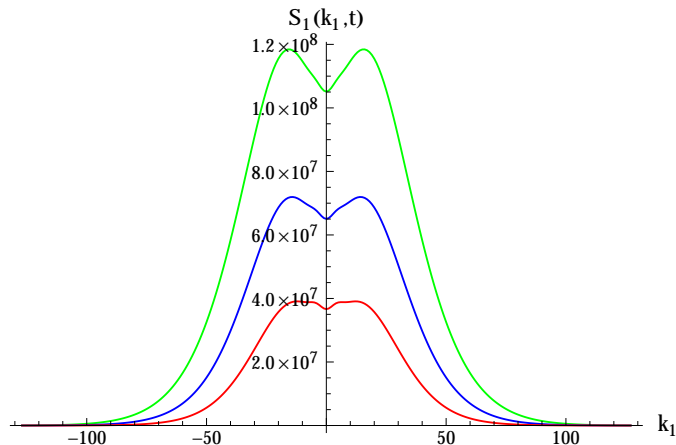


Figure 9: Type II, $a = 20$, $E_0 = 5 \times 10^4$. Plot of the marginal enstrophy density $S_1(k_1, t)$ at $t \cdot 10^7 = 1521, 1544, 1560$. Simulation range $k_3 \in [-19, 2528]$.

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3. Li-Sinai solutions: simulations. Decay rate.

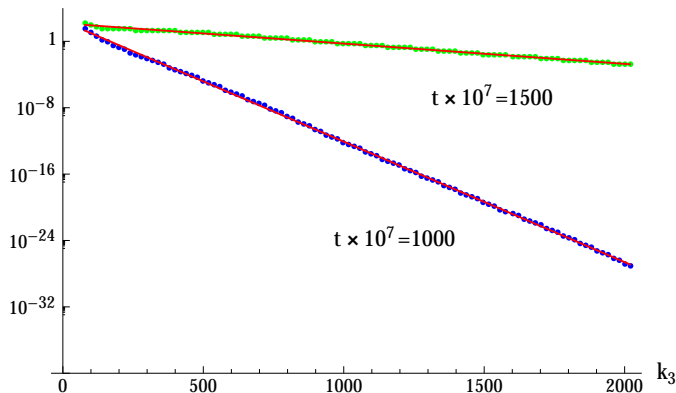


Figure 10: Type II, $a = 20$, $E_0 = 5 \times 10^4$. Plot of $\log(E_3(k_3, t))$, where E_3 is the marginal energy density along the k_3 -axis for $k_3 \geq 400$ at two different times. The dots represent the local maxima of the oscillations of $E_3(k_3, t)$. Simulation range $k_3 \in [-19, 2028]$.

3. Li-Sinai solutions: simulations. Critical time.

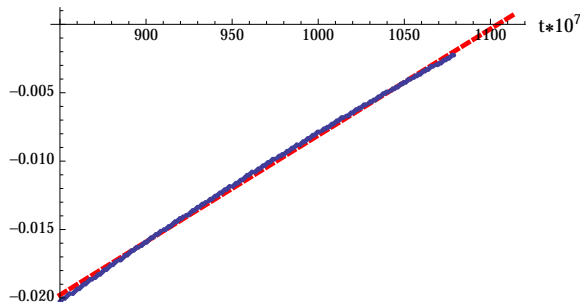


Figure 11: Type I, $a = 20$, $E_0 = 5 \times 10^4$. Exponential decay rate for the marginal density $E_3(k_3, t)$, taken for $k_3 \geq 400$, vs magnified time $t \times 10^7$, with linear regression (dashed line). Simulation range $k_3 \in [-19, 2528]$.

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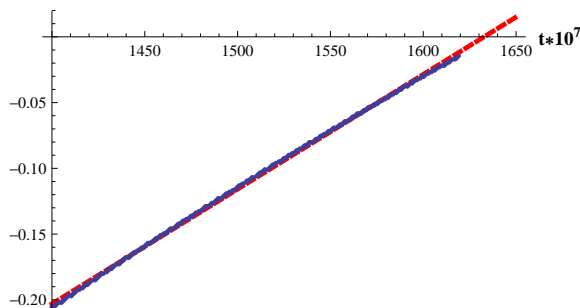


Figure 12: Type II, $a = 20$, $E_0 = 5 \times 10^4$. Exponential decay rate for the marginal density $E_3(k_3, t)$, taken for $k_3 \geq 400$, vs magnified time $t \times 10^7$, with linear regression (dashed line). Simulation range $k_3 \in [-19, 2528]$.

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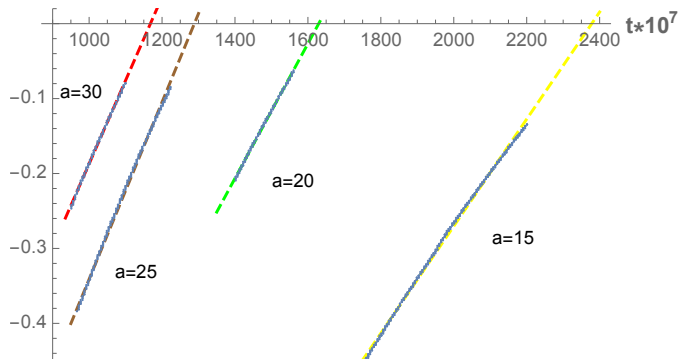


Figure 13: Type II, $E_0 = 5 \times 10^4$. Behavior of the exponential decay rates of $E_3(k_3, t)$ vs. magnified time $t \times 10^7$ for $a = 15, 20, 25, 30$. Simulation range $k_3 \in [-19, 3028]$

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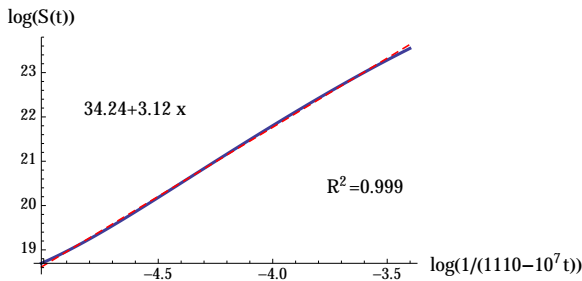


Figure 14: Type I, $a = 20$, $E_0 = 5 \times 10^4$. Log-plot of the total enstrophy $S(t)$ vs $\log \frac{1}{\tau_* - t}$, at times near the blow-up, with linear regression (dashed line, the prediction for the slope is 3.0). Simulation range $k_3 \in [-19, 2528]$.

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$$\mathbf{x}_{\pm}^{(0)} = (0, 0, \pm x_3^{(0)}), \quad x_3^{(0)} \approx \frac{\pi}{a}$$

for the solutions of type II.

3. Li-Sinai solutions: simulations. x-space.

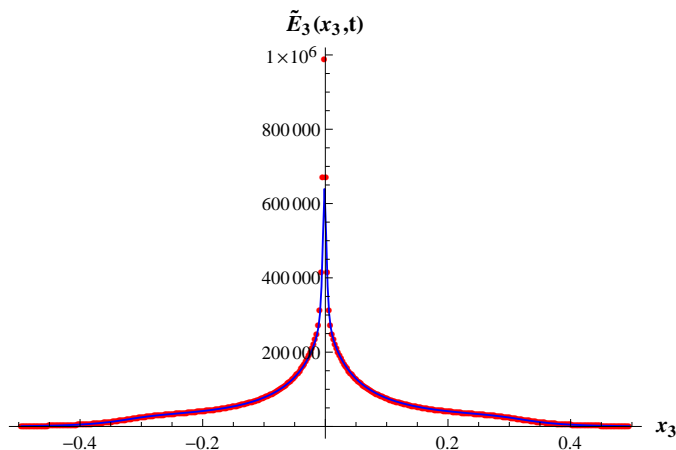


Figure 15: Type I, $a = 20$, $E_0 = 5 \times 10^4$. Plot of the marginal energy density $\tilde{E}_3(x_3, t)$ at $t \cdot 10^7 = 1021$ (dotted line) and $t \cdot 10^7 = 1044$ (continuous line). Simulation range $k_3 \in [-19, 2528]$.

3. Li-Sinai solutions: simulations. x-space.

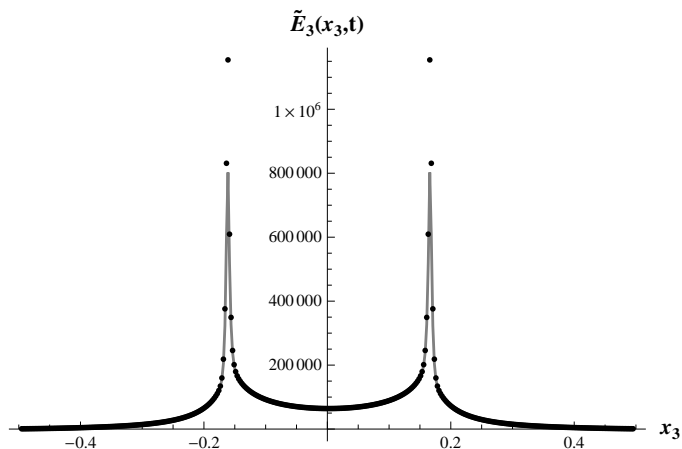


Figure 16: Type II, $a = 20$, $E_0 = 5 \times 10^4$. Plot of the marginal energy density $\tilde{E}_3(x_3, t)$ at $t \cdot 10^7 = 1521$ (continuous line) and $t \cdot 10^7 = 1544$ (dotted line). Simulation range $k_3 \in [-19, 2528]$

3. Li-Sinai solutions: simulations. \mathbf{X} -space.

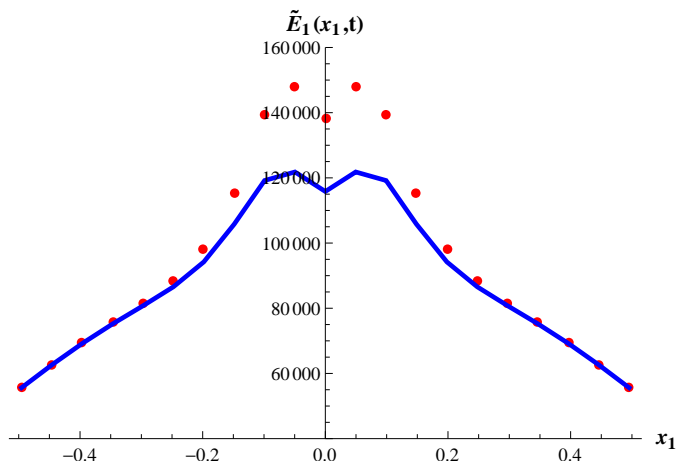


Figure 17: Type II, $a = 20$, $E_0 = 5 \times 10^4$. Plot of the marginal energy density $\tilde{E}_1(x_1, t)$ at $t \cdot 10^7 = 1521$ (continuous line), and $t \cdot 10^7 = 1544$ (dotted line). Simulation range $k_3 \in [-19, 2528]$.

4. REAL SOLUTIONS.

Assuming antisymmetric initial data $A\mathbf{v}_0(\mathbf{k})$

$$\mathbf{v}_0(\mathbf{k}) = \mathbf{v}_0^{\pm}(\mathbf{k}) - \mathbf{v}_0^{\pm}(-\mathbf{k}),$$

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where \mathbf{v}_0^\pm are as above, we get a real solution. (The choice \pm amounts to a change of sign.) $\mathbf{v}_0(\mathbf{k})$ has support in two separate regions around the points $\pm(0, 0, a)$.

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The simulation show, if the initial energy is large enough, that:

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- The total enstrophy $S(t)$ grows, reaches a maximum at a time t_* (which for a fixed depends on E_0), then falls;

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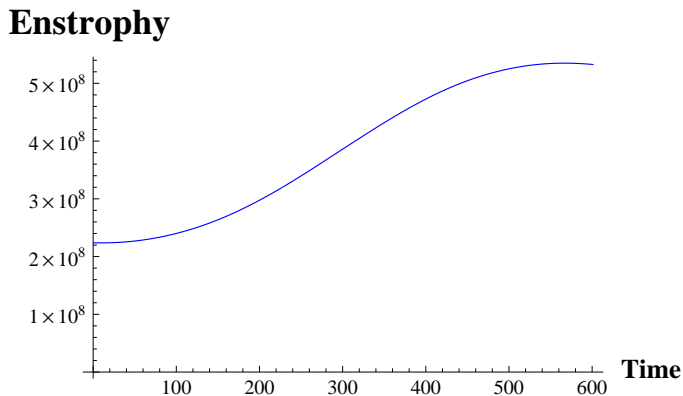


Figure 18: Plot of the total enstrophy $S(t)$ vs. magnified time $t \times 3.2 \times 10^7$. Initial energy $\bar{E}_0 = 2.5 \times 10^4$, $a = 20$.

- The large k_3 modes fall off exponentially fast, with a rate decreasing in absolute value up to $t \approx t_*$, then stays constant.

4. Real solutions.

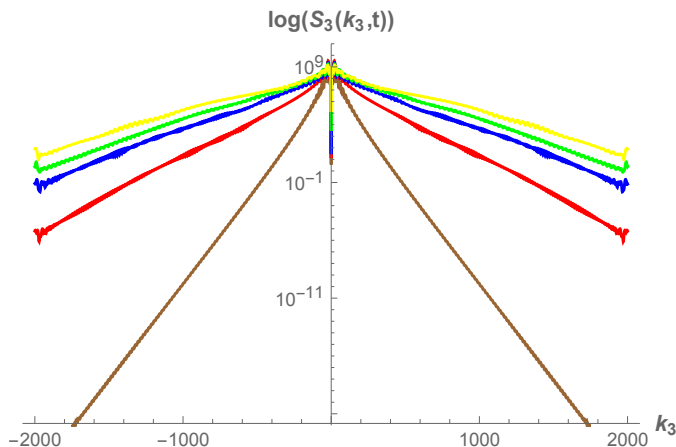


Figure 19: Logarithmic plot of the marginal density $\tilde{E}_3(k_3, t)$ at the (magnified) times $t \times 3.2 \times 10^7 = 100, 200, 300, 400, 500$. Initial energy $\bar{E}_0 = 2.5 \times 10^4$, $a = 20$.

- At the time t_* , the energy and the enstrophy concentrate in two (pseudo)-spikes, close to the singularities of the complex solution of type // with the same a .

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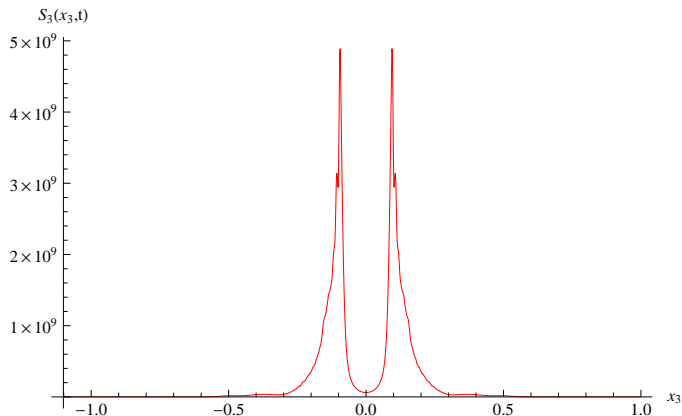


Figure 20: Plot of the marginal density $\tilde{S}_3(x_3, t)$ at $t = 1.27 \times 10^{-5}$. Initial energy $\bar{E}_0 = 2.5 \times 10^4$, $a = 20$.