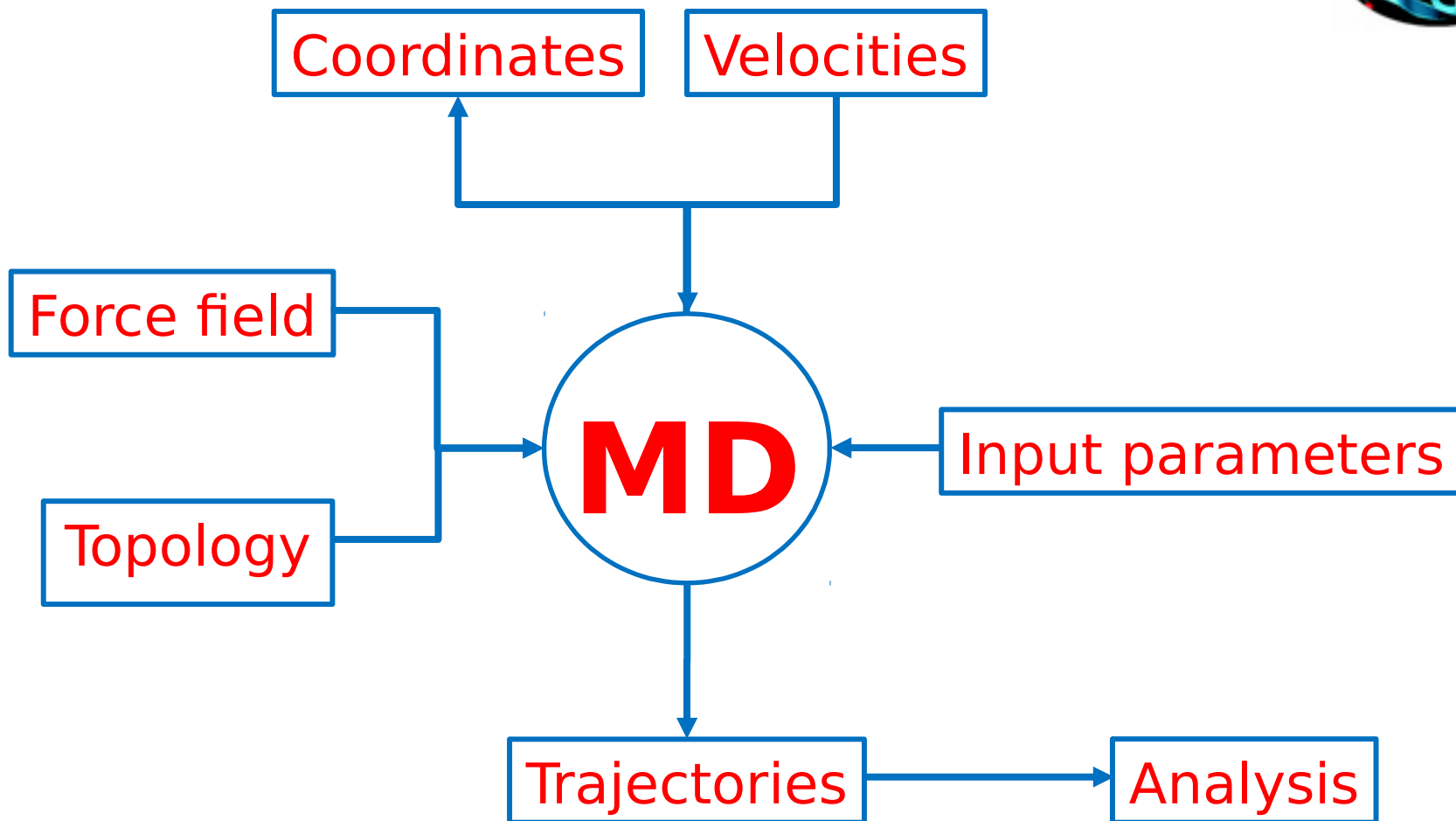
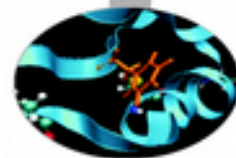


# Introduction to Classical Molecular Dynamics

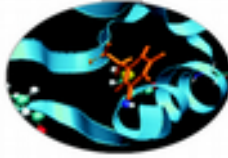


Giovanni Chillemi [g.chillemi@ Cineca.it](mailto:g.chillemi@ Cineca.it)  
HPC department, CINECA

# MD ingredients



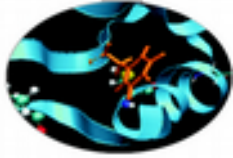
# Equation of motion



The equations that describe the temporal evolution of a physical system is called **equation of motion**. There are different equations of motions, which characterize the motion with different levels of approximation:

- Time-dependent Schrödinger's Equation
  - for quantum-mechanical system
- **Newton's Equation**
  - **for classical-mechanical system**
- Langevin's Equation
  - for stochastic system

# Newton's Equation of motion



Molecules are quantum-mechanical systems whose motion should be described by Schrödinger's Equation. However, technical difficulties make solving Schrödinger's Equation for large systems impractical.

Therefore the motion of a molecule is usually approximated by the laws of classical mechanics and by Newton's equation of Motion. In its most simplistic form Newton's second law of motion states:

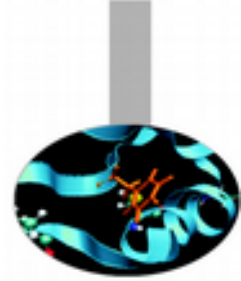
$$\mathbf{f}_i = m_i \cdot \mathbf{a}_i$$

where  $m_i$  is the mass of particle  $i$ ,  $a_i$  is its acceleration. The force  $f_i$  is given as the derivative of the potential energy function  $V$ :

$$\mathbf{f}_i = -\frac{\partial V}{\partial \mathbf{r}_i}$$

where  $r_i$  is the position of particle  $i$

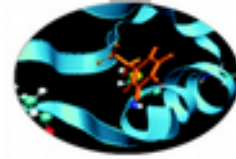
# Potential energy function



$$\begin{aligned}
 V(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n) = & \sum_{bond} \frac{1}{2} k_{b_n} (b_n - b_{0_n})^2 + \sum_{angle} \frac{1}{2} k_{\theta_n} (\theta_n - \theta_{0_n})^2 + \\
 & + \sum_{\substack{improper \\ dihedral}} \frac{1}{2} k_{\xi_n} (\xi_n - \xi_{0_n})^2 + \sum_{dihedral} k_{\phi_n} [1 + \cos(m_n \phi_n - \delta_n)] + \\
 & + \sum_{\substack{nonbonded \\ pairs(ij)}} \left( \left( \frac{C_{ij}^{(12)}}{r_{ij}^{12}} - \frac{C_{ij}^{(6)}}{r_{ij}^6} \right) + \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{\epsilon_r r_{ij}} \right)
 \end{aligned}$$

non bonded interactions

bonded interactions

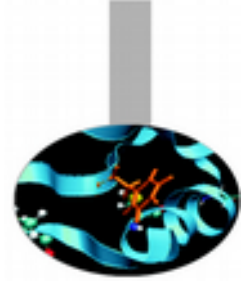


Model	Degree of freedom	Example of predicted properties	
Quantum mechanics	<b>Considered</b> Nucleus, electrons	<b>Removed</b> Nucleons	Chemistry reaction
Polarizable atoms	Atoms, dipoles	Electrons	Binding of charged substrates
<b>Non polarizable atoms</b>	Solute atoms, solvent atoms	Dipoles	Conformational transitions Hydration
Implicit solvent	Solute atoms	Solvent atoms	Folding topology of macromolecules



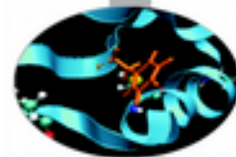
Classical Molecular Dynamics

# A Brief History



Year	System
1964	Liquid Argon (Rahman Phys Rev)
1974	Water (Rahman J. Chem Phys. )
1977	Small protein in vacuo (Mc Cammon Karplus Nature)
1988	First Protein in explicit water (Levitt PNAS )
From 1995	Protein-DNA Complexes – Membrane Proteins- Complex Systems

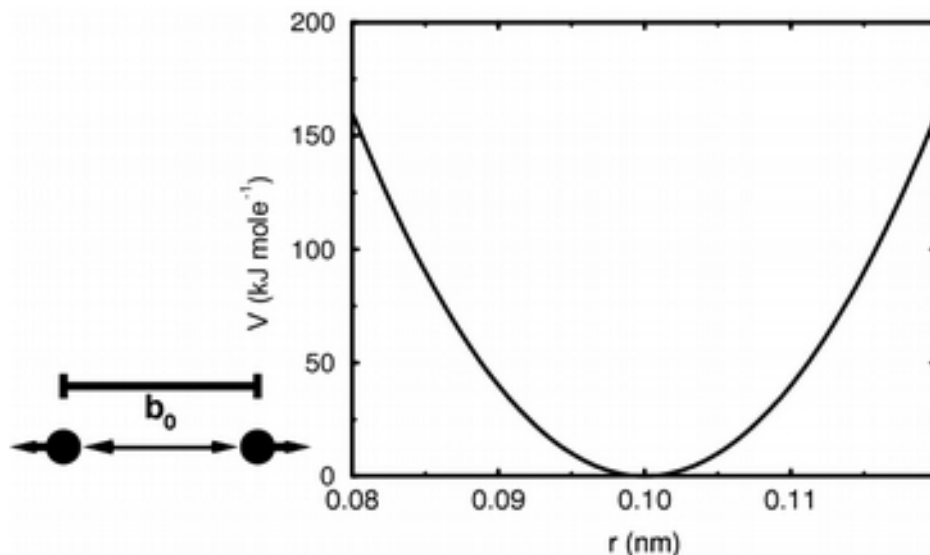
# Bond Stretching Energy



$$= \sum_{\text{bond}} \frac{1}{2} k_{b_n} (b_n - b_{0_n})^2 + ..$$

$k_b$  is the spring constant of the bond  $b_0$  is the bond length at equilibrium

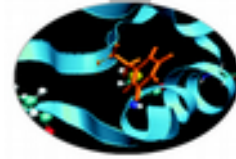
Unique  $k_b$  and  $b_0$  assigned for each bond pair, i.e. C-C, O-H



Principle of bond stretching (left), and the bond stretching potential (right).



# Bond Stretching Force



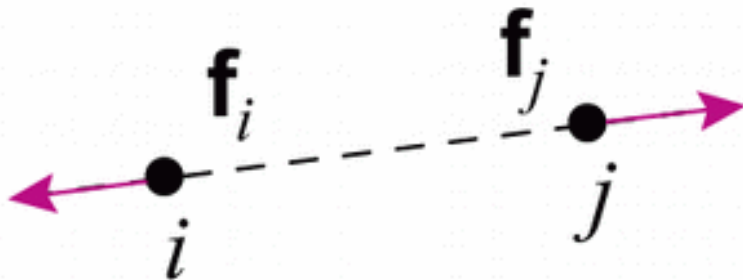
$$\mathbf{f}_i = -\frac{\partial V^{bond}}{\partial \mathbf{r}_i} = -\frac{\partial V^{bond}}{\partial r_{ij}} \frac{\partial r_{ij}}{\partial \mathbf{r}_i} = k_b (r_{ij} - b_0) \frac{\mathbf{r}_{ij}}{r_{ij}}$$

$$\mathbf{f}_j = -\mathbf{f}_i$$

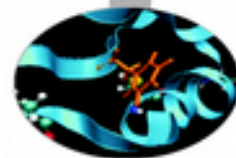


If atom *i* and *j* are closer than  $b_0$ , the bond force separates them

If atom *i* and *j* are farther than  $b_0$ , the bond force draws them nearer

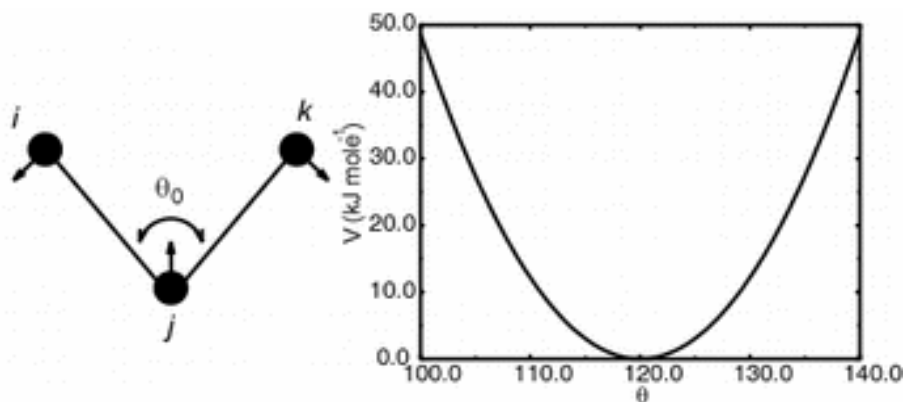


# Bending Energy



$$+ \sum_{\text{angle}} \frac{1}{2} k_{\theta_n} (\theta_n - \theta_{0_n})^2 +$$

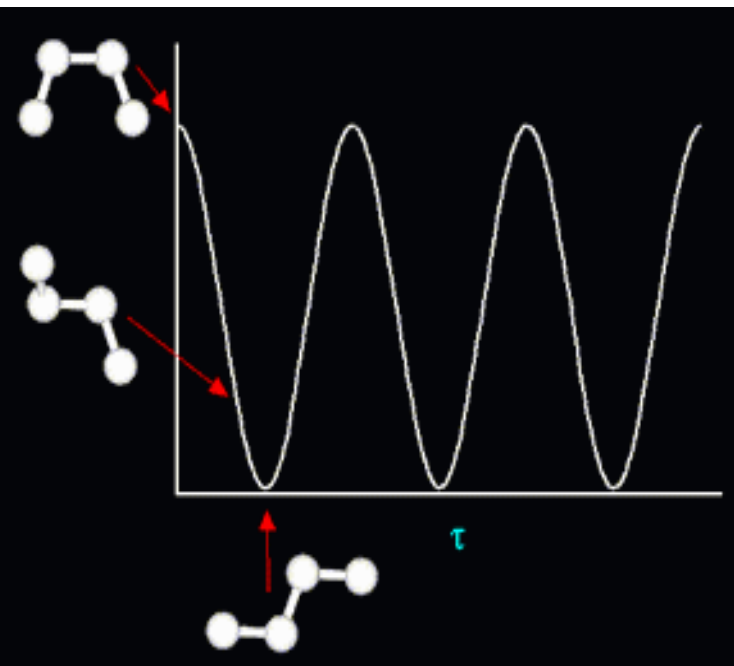
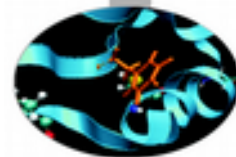
$k_{\theta}$  is the spring constant of the bending  
 $\theta_0$  is the angle bending at equilibrium



Principle of angle vibration (left) and the bond angle potential (right).

Unique parameters for angle bending are assigned to each bonded triplet of atoms based on their types (e.g. C-C-C, C-O-C, C-C-H, etc.)

# Torsional or Dihedral Energy

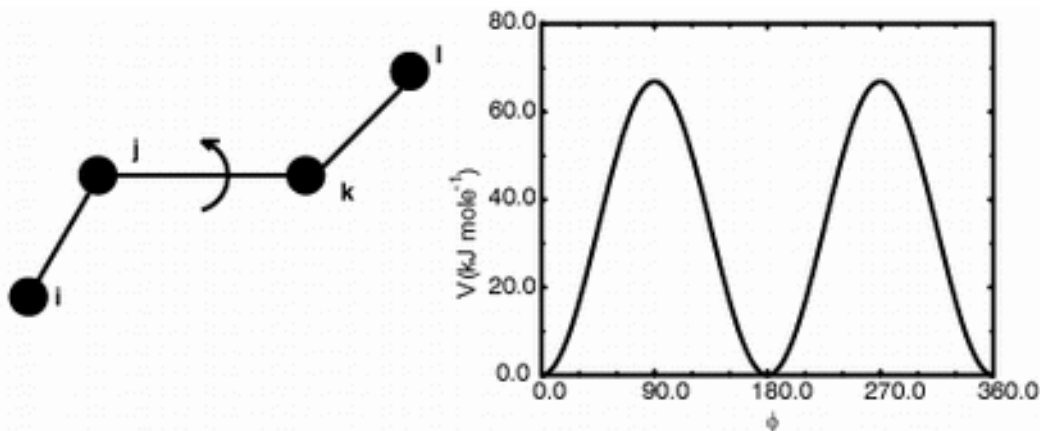


$$+ \sum_{\text{dihedral}} k_{\phi_n} [1 + \cos(m_n \phi_n - \delta_n)] +$$

$\phi$  = angle

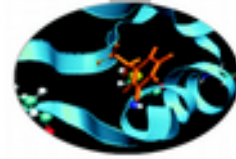
$\delta$  = phase

$m$  = number of peaks in a full rotation

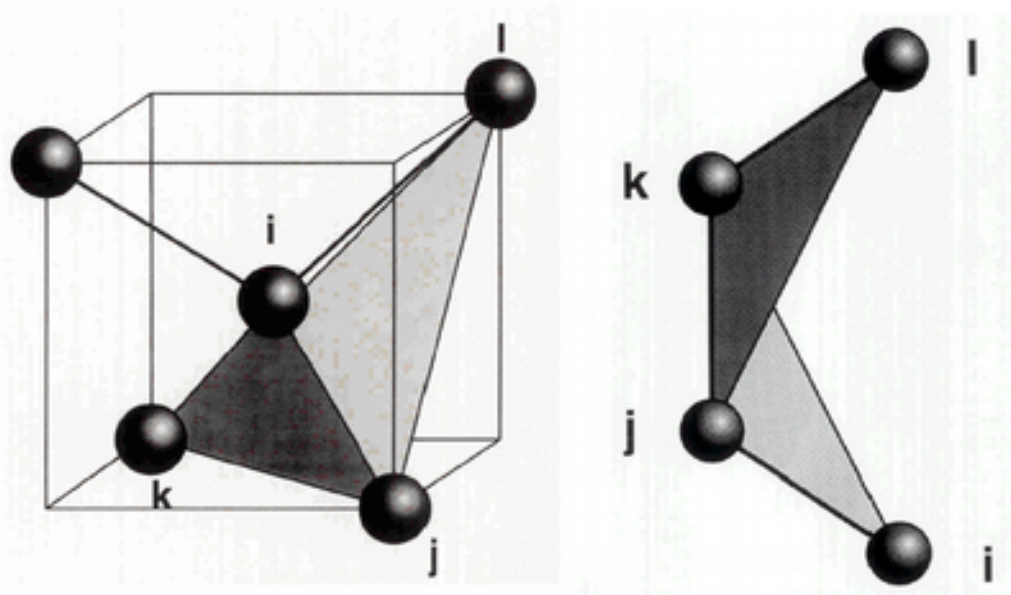


Principle of proper dihedral angle (left, in *trans* form) and the dihedral angle potential (right).

# Improper Dihedral Energy



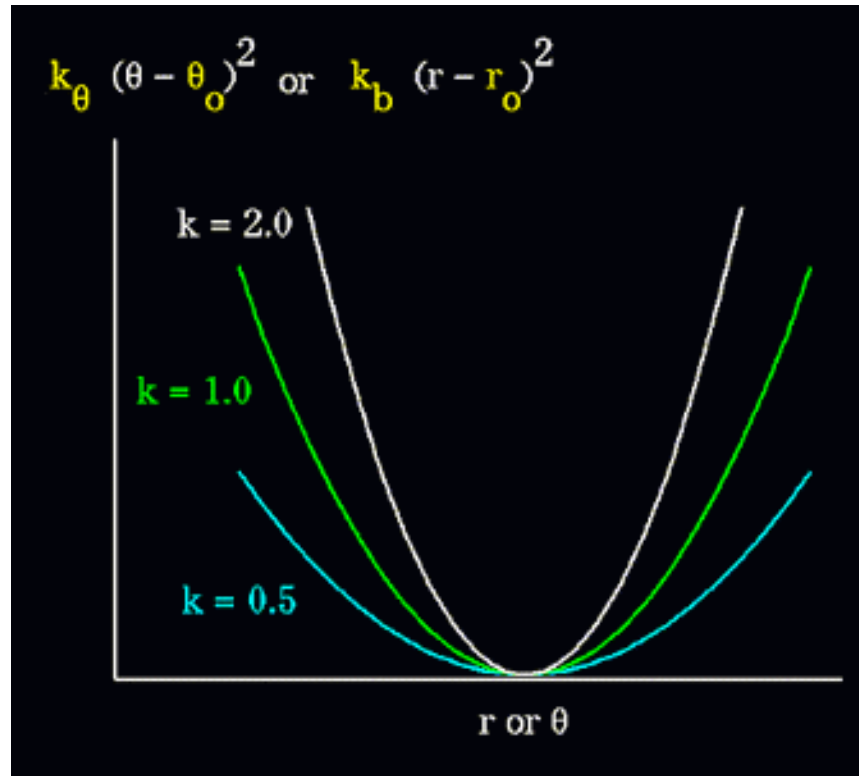
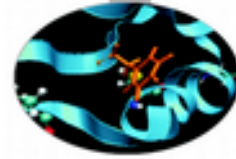
The energy required to deform a group of atoms from its equilibrium angle,  $x_0$ . Used for tetrahedral or planar groups



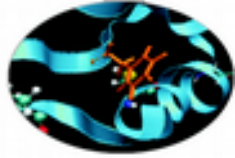
$$+ \sum_{\text{improper dihedral}} \frac{1}{2} k_{\xi_n} (\xi_n - \xi_{0n})^2 +$$

Again this system can be modeled by a spring, and the energy is given by the Hookean potential with respect to the planar angle

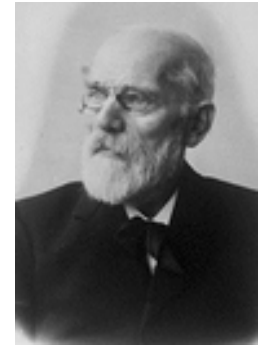
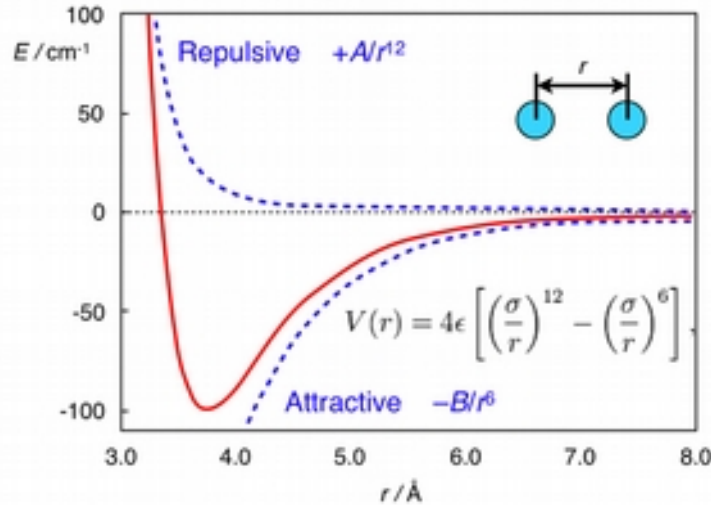
# The “Hookean” potential



$k_b$  and  $k_{\theta}$  broaden or steepen the slope of the parabola  
The larger the value of  $k$ , the more energy is required to deform an angle (or bond) from its equilibrium value



Sir John Lennard Jones



Johannes Diderik Van der Waals

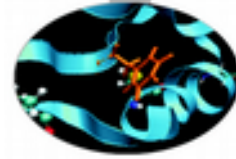
## LJ interactions

$$+ \sum_{\text{nonbonded pairs}(ij)} \left( \left( \frac{C_{ij}^{(12)}}{r_{ij}^{12}} - \frac{C_{ij}^{(6)}}{r_{ij}^6} \right) + \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{\epsilon_r r_{ij}} \right)$$

non bonded interactions

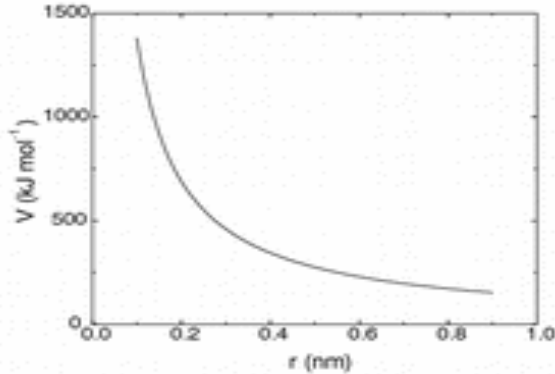


# Electrostatic interactions



$q_i$  and  $q_j$  are the partial atomic charges for atoms  $i$  and  $j$ ,  
separated by a distance  $r_{ij}$

$\epsilon_r$  is the relative dielectric constant



The Coulomb interaction (for particles with equal signed charge)

**Electrostatic**

$$+ \sum_{\text{nonbonded pairs}(ij)} \left( \left( \frac{C_{ij}^{(12)}}{r_{ij}^{12}} - \frac{C_{ij}^{(6)}}{r_{ij}^6} \right) + \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{\epsilon_r r_{ij}} \right)$$

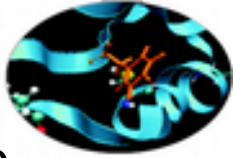
non bonded interactions



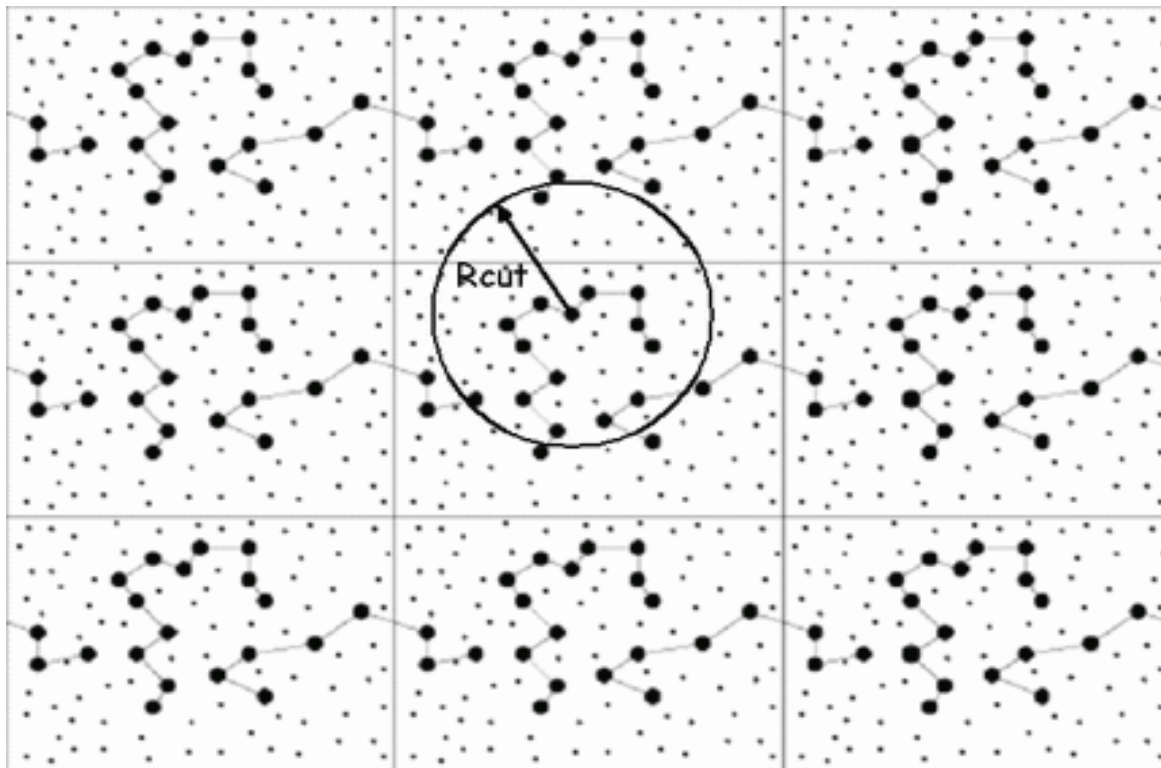
Charles Augustin de Coulomb



# Periodic boundary conditions and cut-off radius



To simulate our finite system in liquid conditions, we apply the pbc: i.e. the system box is virtually surrounded in all directions by copy of itself. An atom close to a box border interacts with the atoms in another pbc image. The non-bonded interactions are only calculated between atom pairs closer than a spherical cut-off

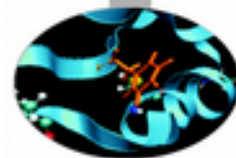




# BOX dimension

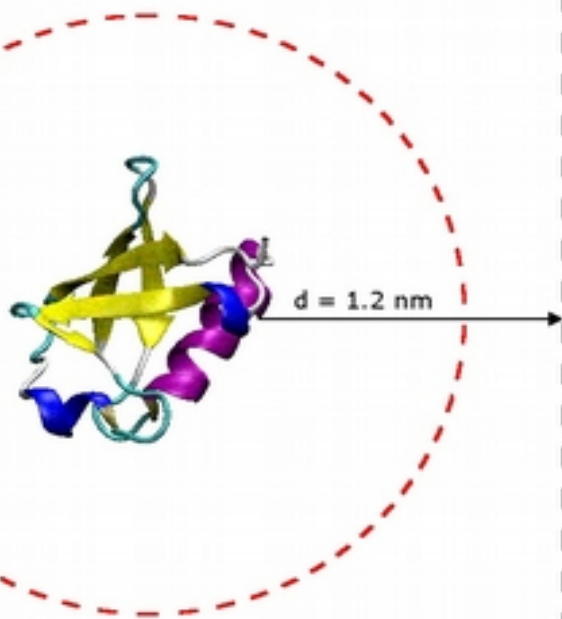
The edge of cubic box must be large enough to avoid interactions of the solute with itself.

Its minimal dimension therefore depends by the chosen cut-off for the non bonded interactions

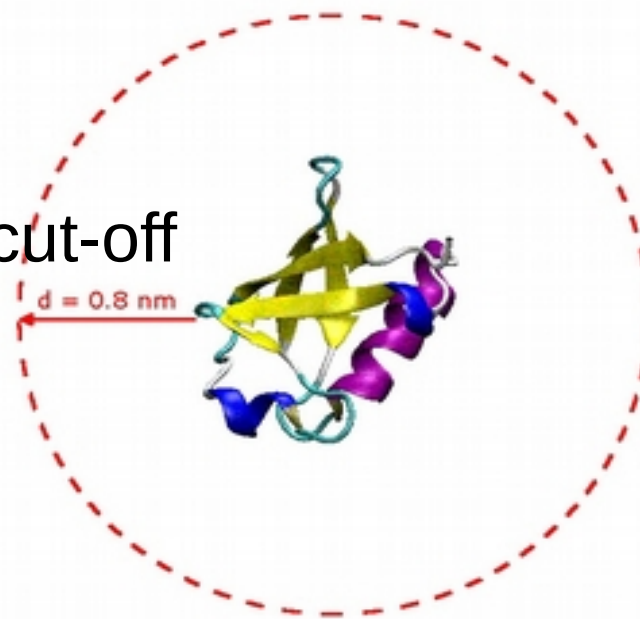


Original Protein

Periodic Image

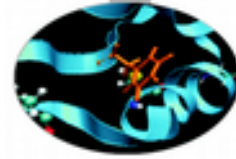


cut-off



Edge of the box = 3.2 nm  
(2.4 nm of water plus the solute size)

# Electrostatic interactions: Particle Mesh Ewald (PME)



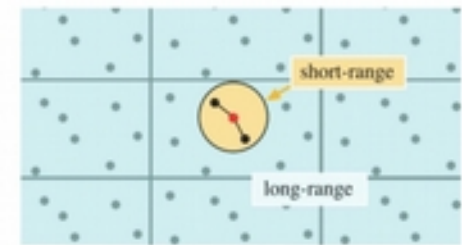
- Short range in the real space
- Long range in the Fourier space

Coulomb forces on  $N$  particles, charges  $q_i$ , positions  $\mathbf{r}_i$ , box length  $L$ , periodic b.c.

- electrostatic potential

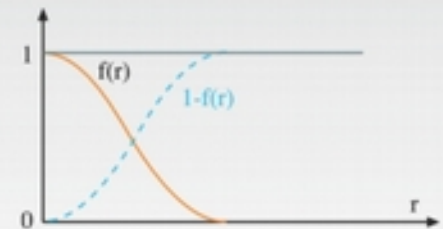
$$V = \frac{1}{2} \sum_{i,j=1}^N \sum_{\mathbf{n} \in \mathbb{Z}^3} \frac{q_i q_j}{|\mathbf{r}_{ij} + \mathbf{n}L|}$$

straightforward summation  
impracticable



- Trick 1: Split problem into 2 parts with help of:

$$\frac{1}{r} = \underbrace{\frac{f(r)}{r}}_{\text{short range}} + \underbrace{\frac{1-f(r)}{r}}_{\text{long range}}$$



$$V = \underbrace{V_{dir}}_{\text{direct space}} + \underbrace{V_{rec}}_{\text{fourier space}}$$

The cut-off radius method for electrostatic interactions is particularly inaccurate for charged molecules such as DNA or for dipolar groups such as alpha helices

PME corrects these errors and it helps maintaining short the cut-off in the real space: i.e. the number of atom pairs is reduced

# Charge groups and atom types

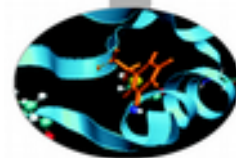
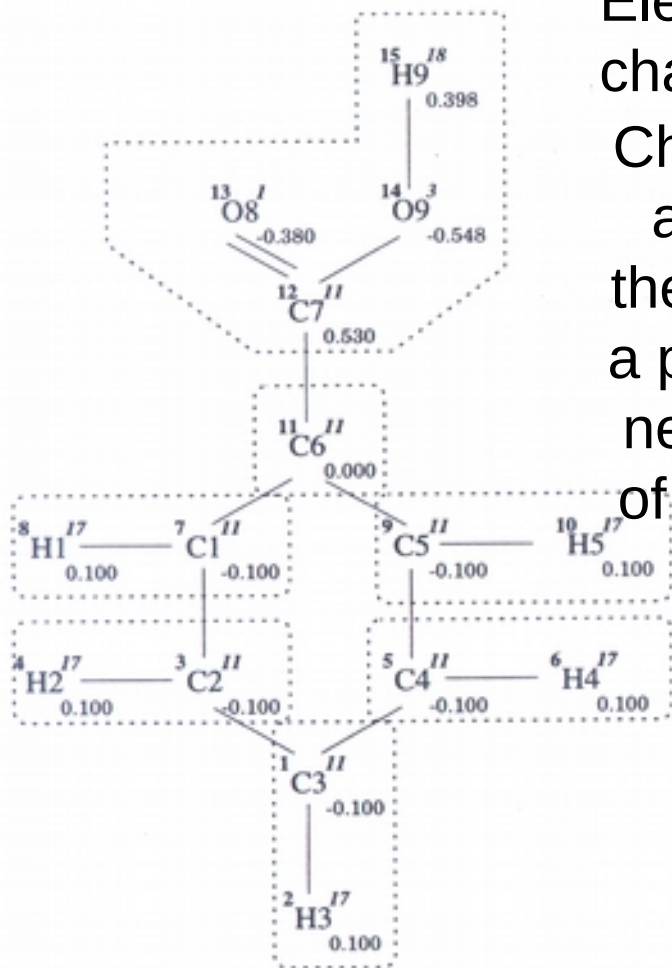


Fig. 4.2.70 Molecular topology building block definition

Solute building block: Benzoic acid (neutral)  
 Name: BA

a. Atoms



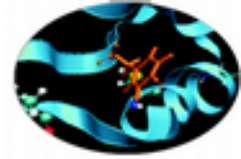
A **charge group** is a neutral charge group composed by several partially charged atoms of a chemical group.

Electrostatics can be calculated between charge groups instead that atom pairs

Charge groups were first introduced to reduce artifacts in the electrostatics calculation but they can also speed up the calculations; given a pair of water molecules for instance, we only need to determine one atom distance instead of nine (or sixteen for a four-site water model)

Note that an **atom type** is not a physical feature. O8 is defined with a different atom type than O9. In fact, their bond constants with C7 and atomic charges are different

# Force field

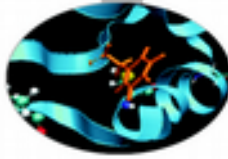


$$\begin{aligned}
 V(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n) = & \sum_{bond} \frac{1}{2} k_{b_n} (b_n - b_{0_n})^2 + \sum_{angle} \frac{1}{2} k_{\theta_n} (\theta_n - \theta_{0_n})^2 + \\
 & + \sum_{improper\ dihedral} \frac{1}{2} k_{\xi_n} (\xi_n - \xi_{0_n})^2 + \sum_{dihedral} k_{\phi_n} [1 + \cos(m_n \phi_n - \delta_n)] + \\
 & + \sum_{nonbonded\ pairs(ij)} \left( \left( \frac{C_{ij}^{(12)}}{r_{ij}^{12}} - \frac{C_{ij}^{(6)}}{r_{ij}^6} \right) + \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{\epsilon_r r_{ij}} \right)
 \end{aligned}$$

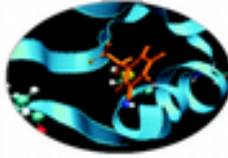
The potential energy function, together with the parameters required to describe the behavior of different kinds of atoms and bonds ( $k_b$ ,  $k_\theta$ ,  $k_\xi$ ,  $C_{ij}$ , ...), is called a **force field**.

Several force fields are currently used and the choice depends from the studied system. Some force field are better suited for nucleic acids, for example, while others for membrane proteins

# Available forcefield in Gromacs (4.6.5)



1. AMBER03 protein, nucleic AMBER94 (Duan et al., J. Comp. Chem. 24, 1999-2012, 2003)
2. AMBER94 force field (Cornell et al., JACS 117, 5179-5197, 1995)
3. AMBER96 protein, nucleic AMBER94 (Kollman et al., Acc. Chem. Res. 29, 461-469, 1996)
4. AMBER99 protein, nucleic AMBER94 (Wang et al., J. Comp. Chem. 21, 1049-1074, 2000)
5. AMBER99SB protein, nucleic AMBER94 (Hornak et al., Proteins 65, 712-725, 2006)
6. AMBER99SB-ILDN protein, nucleic AMBER94 (Lindorff-Larsen et al., Proteins 78, 1950-58, 2010)
7. AMBERGS force field (Garcia & Sanbonmatsu, PNAS 99, 2782-2787, 2002)
8. CHARMM27 all-atom force field (with CMAP) - version 2.0
9. GROMOS96 43a1 force field
10. GROMOS96 43a2 force field (improved alkane dihedrals)
11. GROMOS96 45a3 force field (Schuler JCC 2001 22 1205)
12. GROMOS96 53a5 force field (JCC 2004 vol 25 pag 1656)
13. GROMOS96 53a6 force field (JCC 2004 vol 25 pag 1656)
14. GROMOS96 54a7 force field (Eur. Biophys. J. (2011), 40,, 843-856)
15. OPLS-AA/L all-atom force field (2001 aminoacid dihedrals)
16. [DEPRECATED] Encad all-atom force field, using full solvent charges
17. [DEPRECATED] Encad all-atom force field, using scaled-down vacuum charges
18. [DEPRECATED] Gromacs force field (see manual)
19. [DEPRECATED] Gromacs force field with hydrogens for NMR

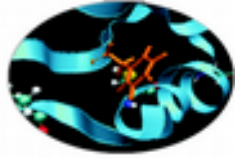


Numeric integration of Newton's equation of motion is typically done step by step using methods that are called **Finite Difference** methods.

These methods use the information available at time  $t$  to predict the system's coordinates and velocities at a time  $t + \delta t$ , where  $\delta t$  is a short time interval and are based on a Taylor expansion of the position at time  $t + \delta t$

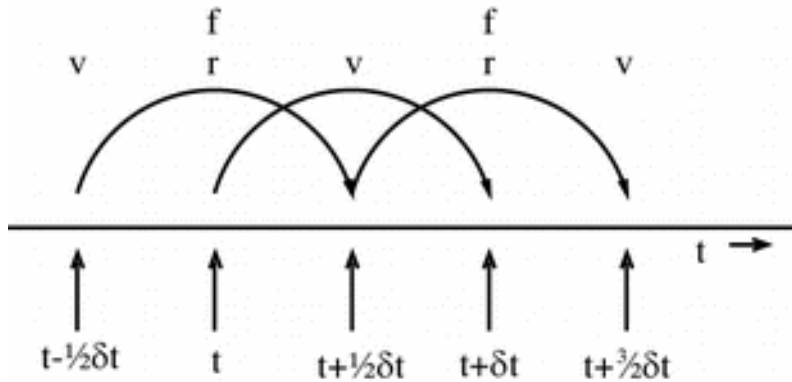
$$r(t + \delta t) = r(t) + v(t)\delta t + \frac{1}{2} a(t)\delta t^2 + \dots$$





$$r(t + \delta t) = 2r(t) - r(t - \delta t) + a(t)\delta t^2$$

❖ *Verlet integrator*



❖ *Leap-frog integrator*

$$r(t + \delta t) = r(t) - v(t + \frac{1}{2} \delta t) \delta t$$

$$v(t + \frac{1}{2} \delta t) = v(t - \frac{1}{2} \delta t) + a(t) \delta t$$

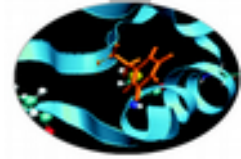
$$r(t + \delta t) = r(t) + v(t)\delta t + \frac{1}{2} a(t)\delta t^2$$

$$v(t + \delta t) = v(t) + [a(t) + a(t + \delta t)] \frac{\delta t}{2}$$

❖ *Velocity Verlet*



# Choice of the timestep



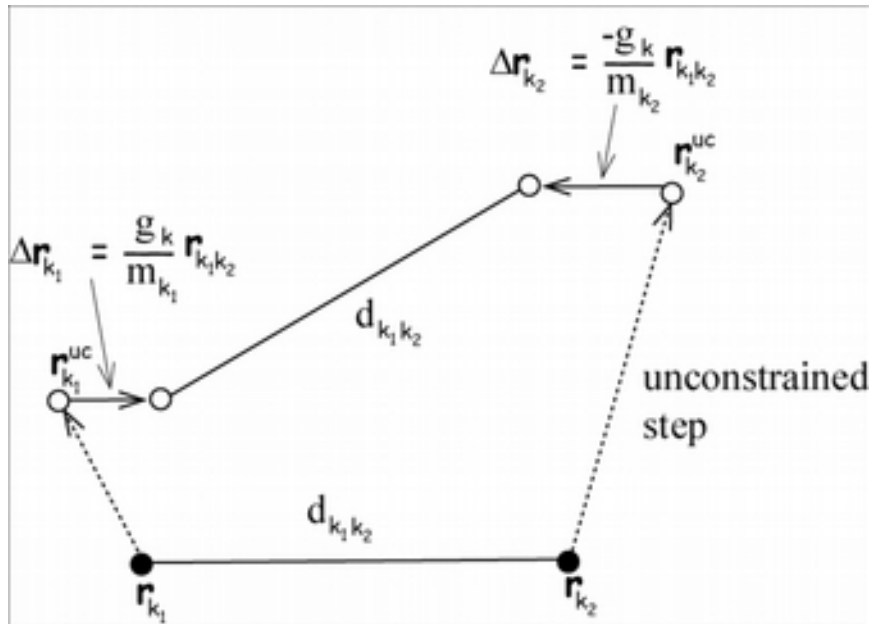
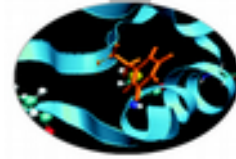
The length of the timestep must be small compared to the period of the highest frequency motions being simulated

Force characteristics	Relaxation time (fs)	Time step (fs)
<b>High frequency motion:</b> bond stretching vibrations	10	0.5
<b>Medium frequency motion:</b> angle bending, proper and improper dihedral angle deformation, LJ and short range Coulombian interactions	40	2
<b>Low frequency motion:</b> long range coulombian interactions	1000	20

The bond stretching vibrations are generally of minimal interest in the study of biomolecular structure and function. Therefore this degree of freedom is usually kept frozen with constraint algorithms ( $2 \times 10^{-15}$  s)



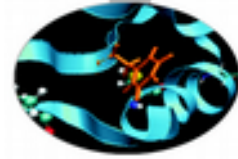
# Constraints



The application of geometrical constraints to maintain fixed all the bond lengths, during a simulation, allows the use of a time step up to 2 fs

In SHAKE (the first constraints algorithm to be implemented in a MD code) changes a set of unconstrained coordinates to a set of coordinates that fulfill a list of distance constraints, solving a set of **Lagrange multipliers** in the constrained equations of motion **(iterative)**

# Constraints

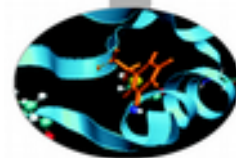


The SETTLE algorithm (developed in 1992) is an analytical solution of SHAKE, specifically for water **(non-iterative)**



Figure 3.9: The three position updates needed for one time step. The dashed line is the old bond of length  $d$ , the solid lines are the new bonds.  $l = d \cos \theta$  and  $p = (2d^2 - l^2)^{\frac{1}{2}}$ .

The LINCS algorithm (developed in 1997, twenty years after SHAKE) solves the bond length constraints, always in two steps **(non-iterative)**



## constraints:

### none

No constraints except for those defined explicitly in the topology, *i.e.* bonds are represented by a harmonic (or other) potential or a Morse potential (depending on the setting of `morse`) and angles by a harmonic (or other) potential.

### h-bonds

Convert the bonds with H-atoms to constraints.

### all-bonds

Convert all bonds to constraints.

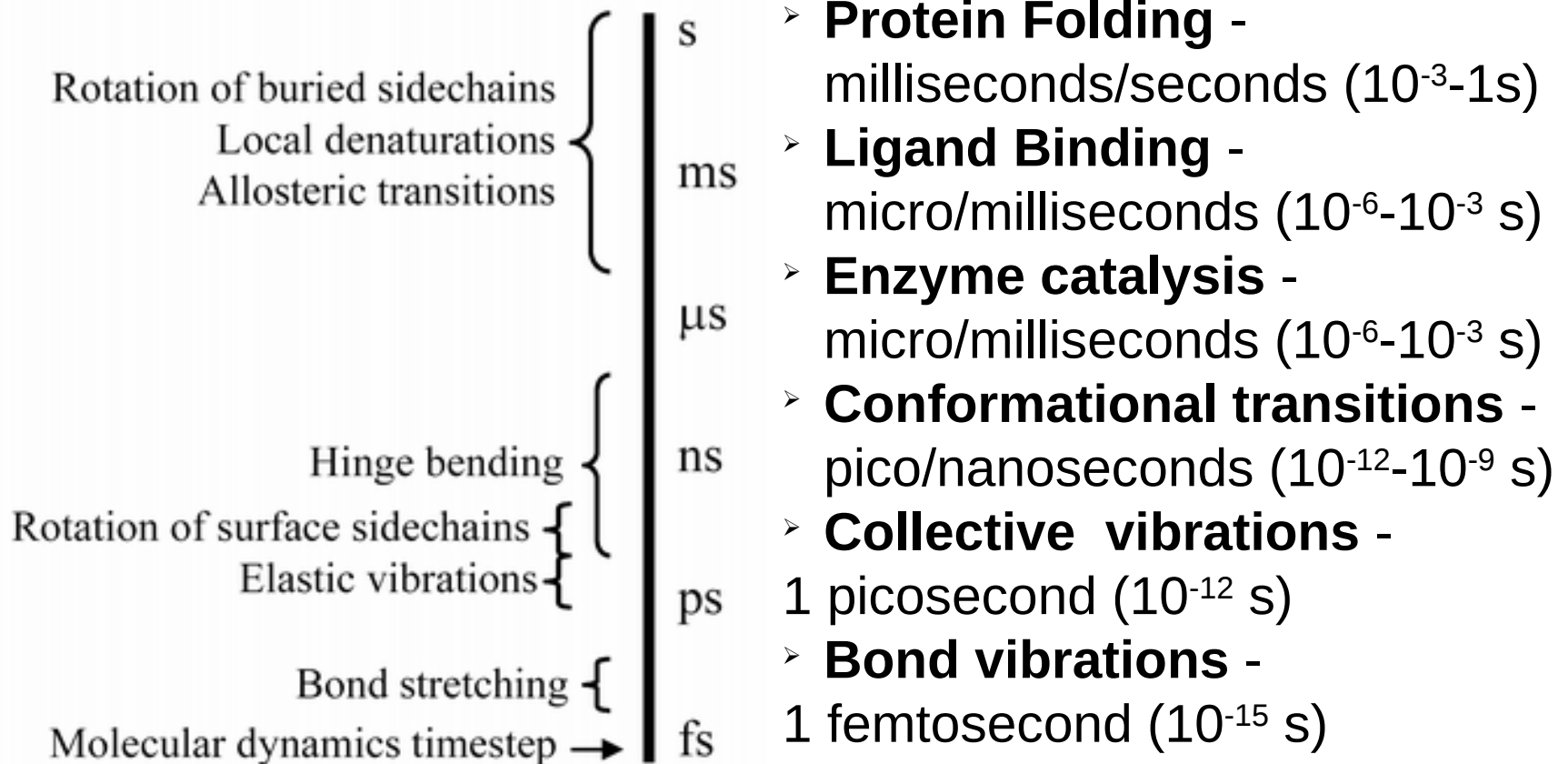
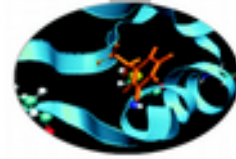
### h-angles

Convert all bonds and additionally the angles that involve H-atoms to bond-constraints.

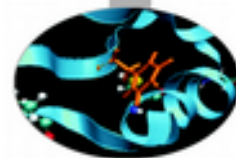
### all-angles

Convert all bonds and angles to bond-constraints.

# Timescale



# Topology



The topology file describes the atoms composing a molecule and their bond connections

Es: flexspc.itp in gromacs

```
[ moleculetype ]
```

```
; molname      nrexcl
SOL            2
```

```
[ atoms ]
```

; id	at type	res nr	res name	at name	cg nr	charge	mass
1	OW_spc	1	SOL	OW	1	-0.82	15.99940
2	HW_spc	1	SOL	HW1	1	0.41	1.00800
3	HW_spc	1	SOL	HW2	1	0.41	1.00800

```
[ bonds ]
```

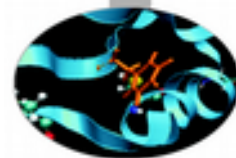
; i	j	funct	length	force.c.
1	2	1	0.1	345000 0.1 345000
1	3	1	0.1	345000 0.1 345000

```
[ angles ]
```

; i	j	k	funct	angle	force.c.
2	1	3	1	109.47 383	109.47 383



# Constraints in Topology



Only in case of water, the constraint algorithm can be selected in the topology file

[ bonds ]

;i	j	funct	length	force.c.
1	2	1	0.1	345000
1	3	1	0.1	345000

flexspc.itp

[ angles ]

;i	j	k	funct	angle	force.c.
2	1	3	1	109	383

[ settles ]

;OW	funct	doh	dhh
1	1	0.1	0.16333

spc.itp

# Topology

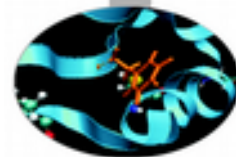
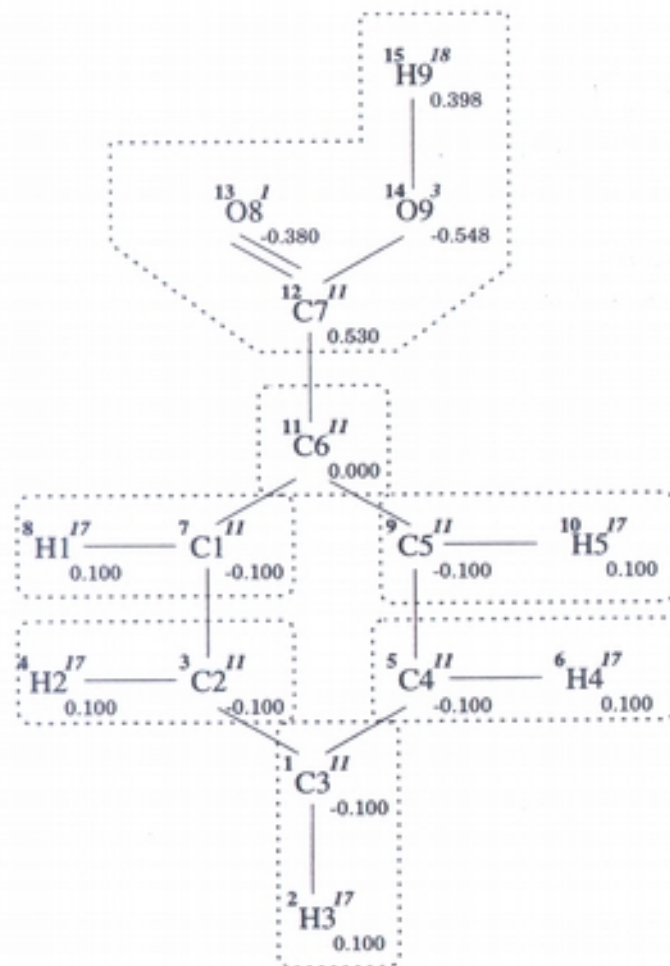


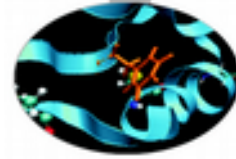
Fig. 4.2.70 Molecular topology building block definition

Solute building block: Benzoic acid (neutral)  
 Name: BA

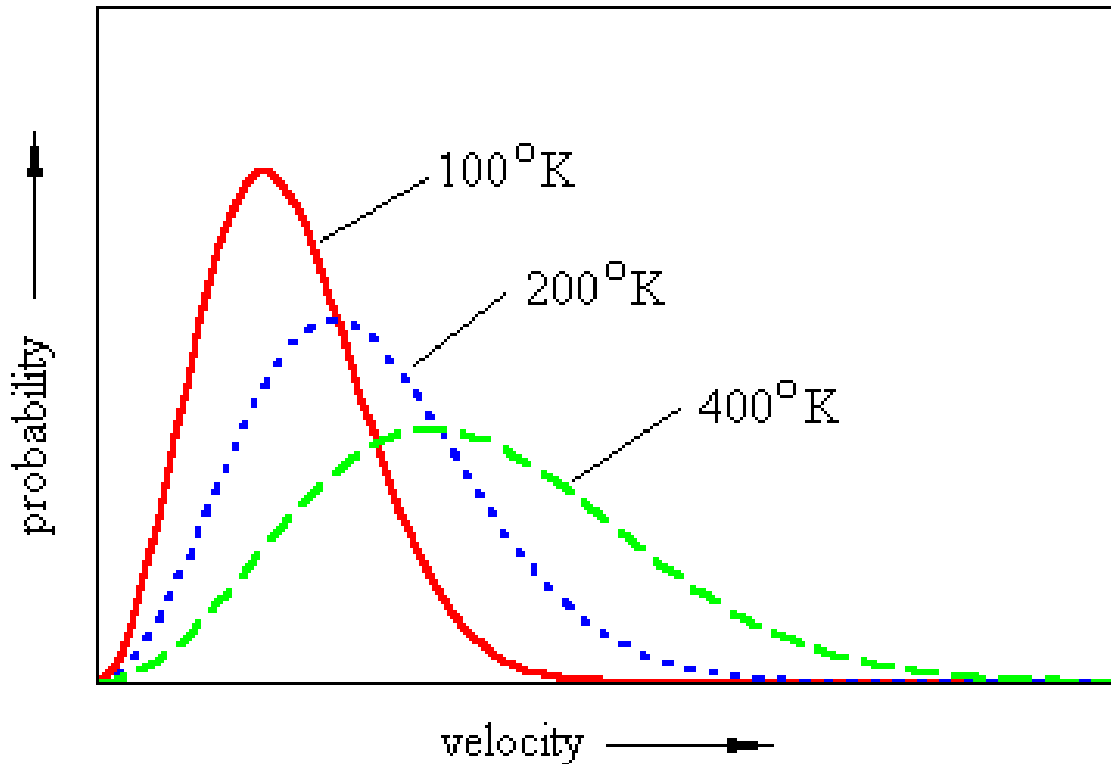
a. Atoms



# Initial velocities

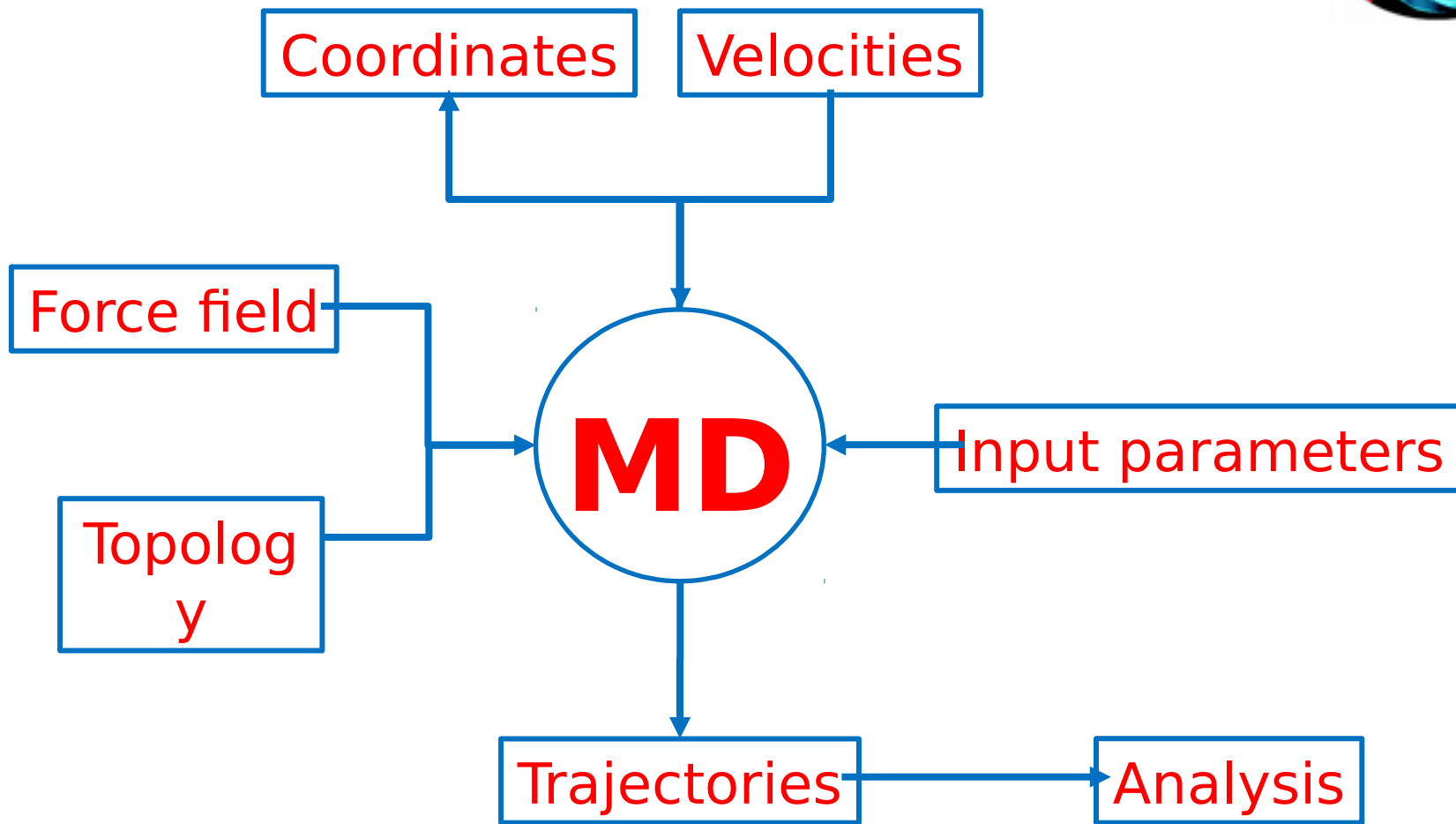
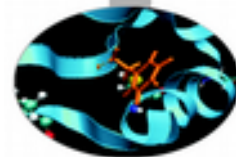


The initial velocity of each atom is random assigned through a **Maxwell-Boltzmann** distribution that is function of the temperature





# To recapitulate..



## THE GLOBAL MD ALGORITHM

### 1. Input initial conditions

Potential interaction  $V$  as a function of atom positions

Positions  $r$  of all atoms in the system

Velocities  $v$  of all atoms in the system



repeat 2,3,4 for the required number of steps:

### 2. Compute forces

The force on any atom

$$F_i = -\frac{\partial V}{\partial r_i}$$

is computed by calculating the force between non-bonded atom pairs:

$$F_i = \sum_j F_{ij}$$

plus the forces due to bonded interactions (which may depend on 1, 2, 3, or 4 atoms), plus restraining and/or external forces.

The potential and kinetic energies and the pressure tensor are computed.



### 3. Update configuration

The movement of the atoms is simulated by numerically solving Newton's equations of motion

$$\frac{d^2 r_i}{dt^2} = \frac{F_i}{m_i}$$

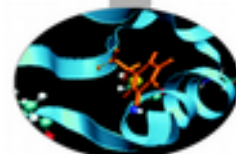
or

$$\frac{dr_i}{dt} = v_i; \quad \frac{dv_i}{dt} = \frac{F_i}{m_i}$$

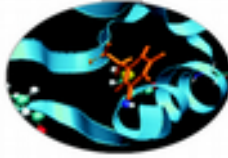


### 4. if required: Output step

write positions, velocities, energies, temperature, pressure, etc.



# Molecular Dynamics ensembles



The method discussed above is appropriate for the micro-canonical ensemble: constant  $N$  (number of particles),  $V$  (volume) and  $E_T$  (total energy =  $E + E_{kin}$ )

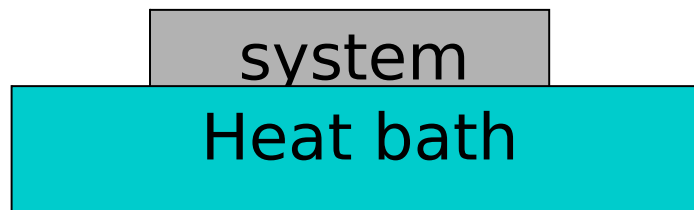
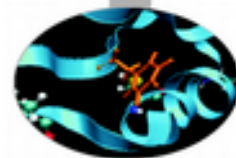
Note that if time step is short enough, the system loses/gains no net energy (potential + kinetic) when running MD in the NVE ensemble

When simulating biological macromolecules, it might be more appropriate to simulate under constant Temperature ( $T$ ) or constant Pressure ( $P$ ):

Canonical ensemble:  $NVT$

Isothermal-isobaric:  $NPT$

# Simulating at constant T: the Berendsen scheme



Bath supplies or removes heat from the system as appropriate

$$\frac{dT(t)}{dt} = \frac{T_0 - T(t)}{\tau_T}$$

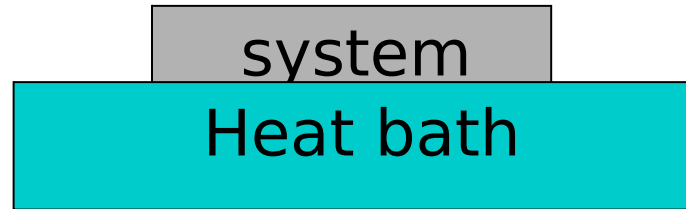
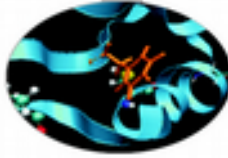
where  $\tau$  determines how strong the bath influences the system

Exponentially scale the velocities at each time step by the factor  $\lambda$ :

$$\lambda = \left[ 1 + \frac{\Delta t}{\tau_T} \left( \frac{T_0}{T(t)} - 1 \right) \right]^{\frac{1}{2}}$$

*T*: “kinetic” temperature

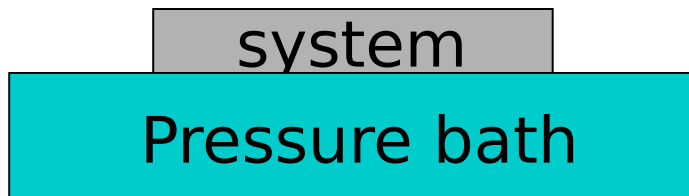
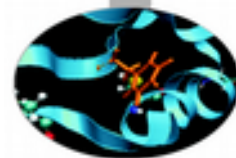
# Simulating at constant T: the Berendsen scheme



A small  $\tau$ , close to the timestep (strong thermostat), is useful in the equilibration phase, when the quick decreasing of the potential energy could increase too much the kinetic energy of the protein

A bigger  $\tau$ , e.g. equal to ten times the timestep (weak thermostat), is useful in the production phase, when we want to keep at minimum the perturbation to the conformational sampling

# Simulating at constant P: the Berendsen scheme



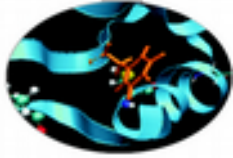
Couple the system to a pressure bath  $\frac{dP(t)}{dt} = \frac{P_0 - P(t)}{\tau_P}$

A change in pressure  $P$  is related to a change in volume  $V$

To regulate pressure: exponentially scale the volume of the simulation box at each time step by a factor  $\mu$

where  $k_T$  : isothermal compressibility  $\mu(t) = \left[ 1 - k_T \frac{\Delta t}{\tau_P} (P_0 - P(t)) \right]^{\frac{1}{3}}$   
 $\tau_P$  : coupling constant

# Sample input file of gromacs



<http://manual.gromacs.org/current/online/mdp.html>

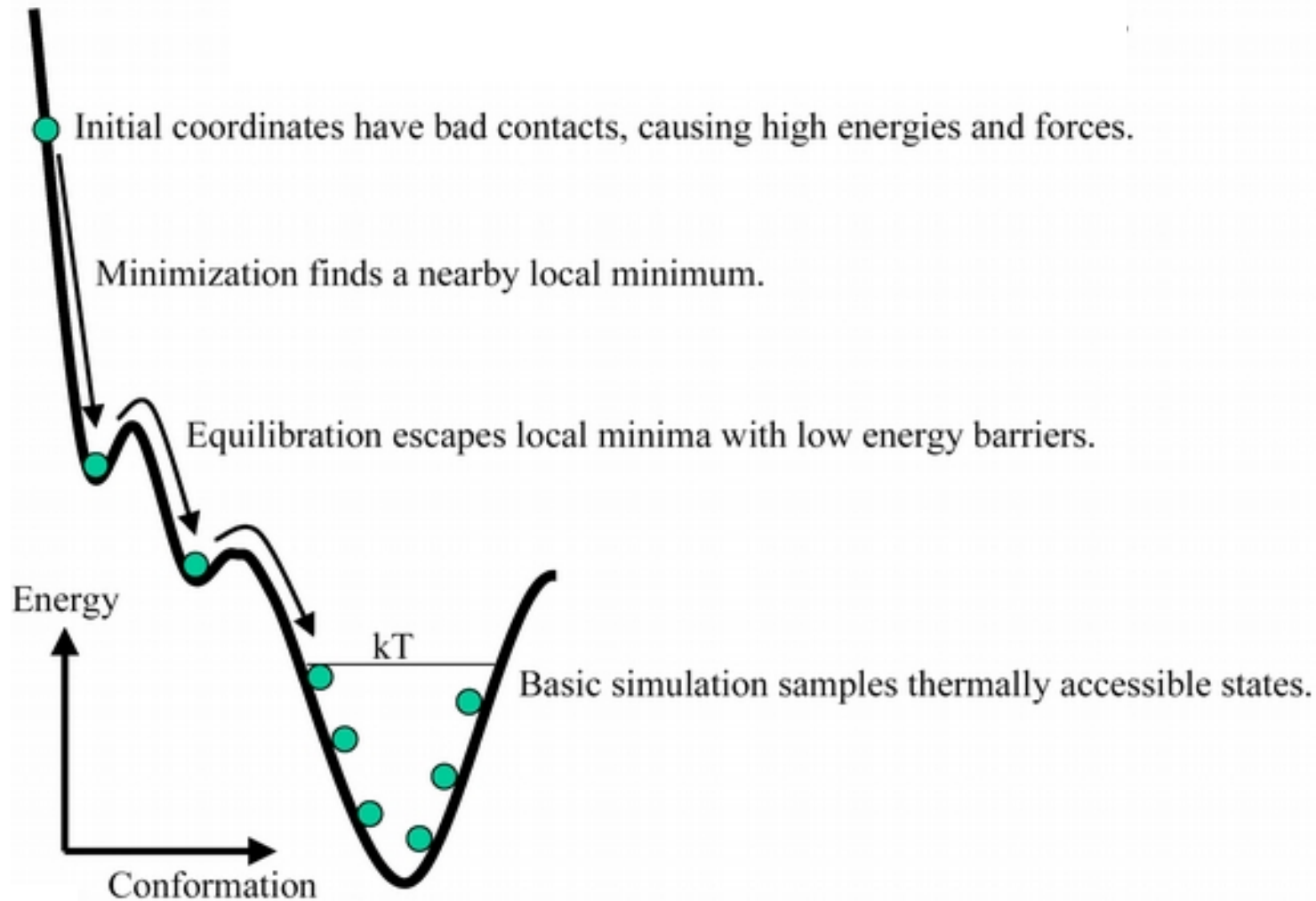
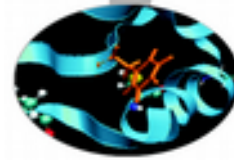
```

title                = Yo
cpp                  = /lib/cpp
include              = -I../top
define               =
integrator           = md
dt                   = 0.002
nsteps               = 500000
nstxout              = 5000
nstvout              = 5000
nstlog               = 5000
nstenergy            = 250
nstxout-compressed   = 250
compressed-x-grps    = Protein
energygrps           = Protein SOL
nstlist              = 10
ns-type              = grid
rlist                = 0.8
coulombtype          = cut-off
rcoulomb             = 1.4
rvdw                 = 0.8

tcoupl               = Berendsen
tc-grps              = Protein SOL
tau-t                = 0.1 0.1
ref-t                = 300 300
Pcoupl               = Berendsen
tau-p                = 1.0
compressibility      = 4.5e-5
ref-p                = 1.0
gen-vel              = yes
gen-temp             = 300
gen-seed             = 173529
constraints          = all-bonds
    
```

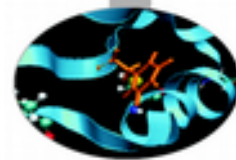


# Conformational sampling

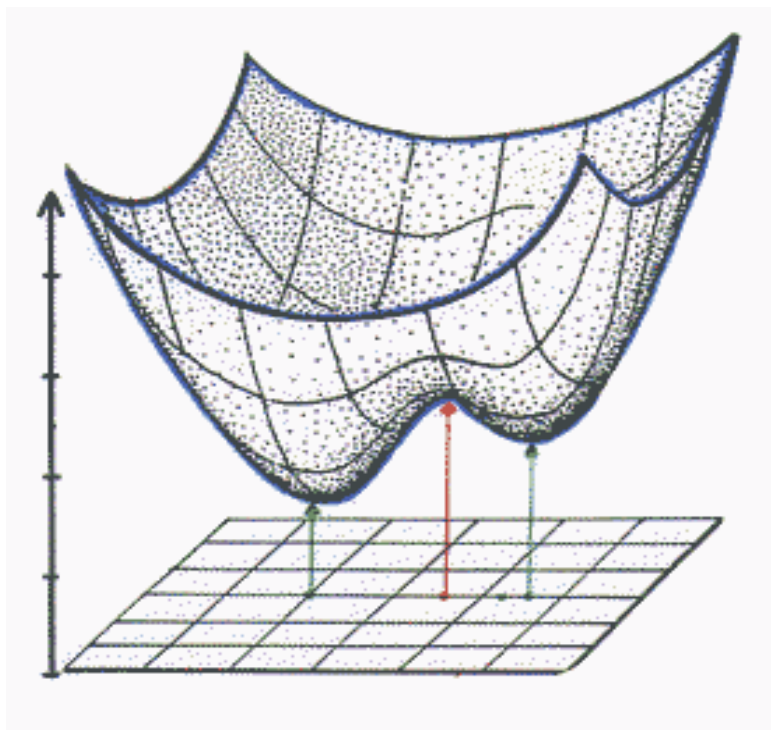


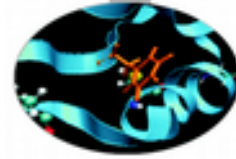


# Energy minimization



The potential energy surface of a molecule is defined by only a global minimum and a great number of local minima: i.e. conformations where all the **first derivative** of the potential energy function with respect to the coordinates are **zero** and all **second derivatives** are **non-negative**





The energy minimization algorithms finds the nearest local minimum, i.e. the minimum that can be reached by systematically moving down the steepest local gradient

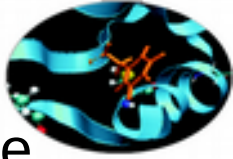
Usually they cannot found the global minumum

Two algorithms very used in MD codes are

- *steepest descent*
- *conjugate gradient*

Both of the first order: they use the first derivative of the energy potential function with respect to the coordinates

The *steepest descent* uses only the gradient of the potential energy function in the local position to calculate the coordinate movement

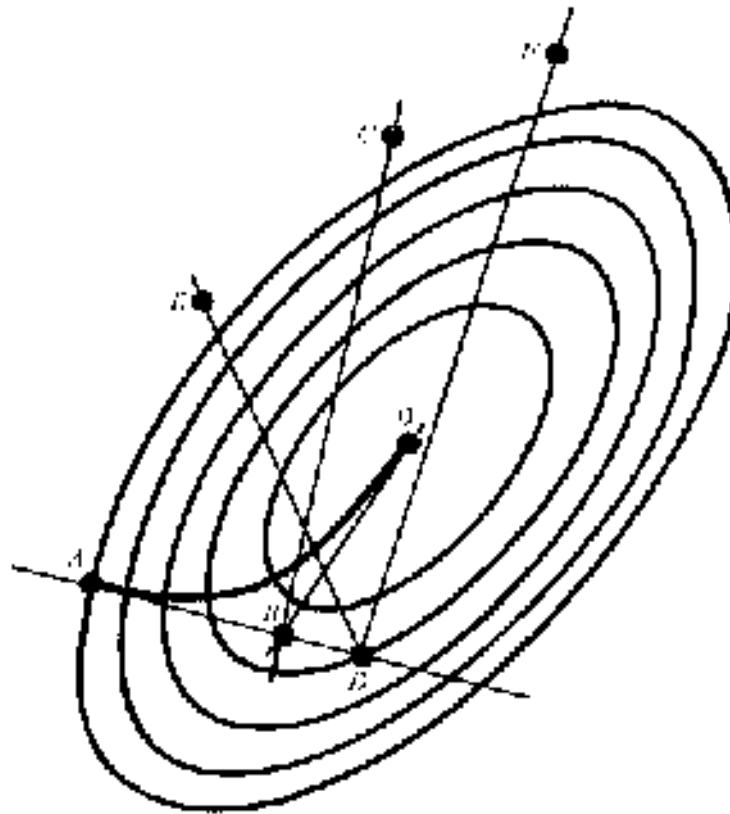
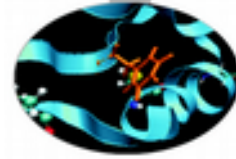


It is quicker in the single iteration but less precise in finding the local minimum. Therefore is useful in the first steps of minimizations

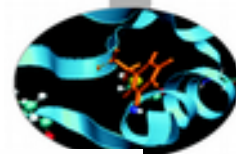
The *conjugate gradient* uses also the gradient of the potential energy function in the previous step

It is more accurate in finding the local minimum but it is slower than the steepest descent. Therefore is usually applied after some steps of steepest descent





Starting from A, the *steepest descents* goes through A-B-C (or A-D-F)  
While the *conjugate gradient*, weighting the A-B and B-C gradients,  
goes through A-B-O



## steep

A steepest descent algorithm for energy minimization. The maximum step size is **emstep** [nm], the tolerance is **emtol** [kJ mol<sup>-1</sup> nm<sup>-1</sup>].

## cg

A conjugate gradient algorithm for energy minimization, the tolerance is **emtol** [kJ mol<sup>-1</sup> nm<sup>-1</sup>]. CG is more efficient when a steepest descent step is done every once in a while, this is determined by **nstcgsteep**.

**emtol: (100.0) [kJ mol<sup>-1</sup> nm<sup>-1</sup>]**

the minimization is converged when the maximum force is smaller than this value

**emstep: (0.01) [nm]**

initial step-size

**nstcgsteep: (1000) [steps]**

frequency of performing 1 steepest descent step while doing conjugate gradient energy minimization.

# 1-4 interactions

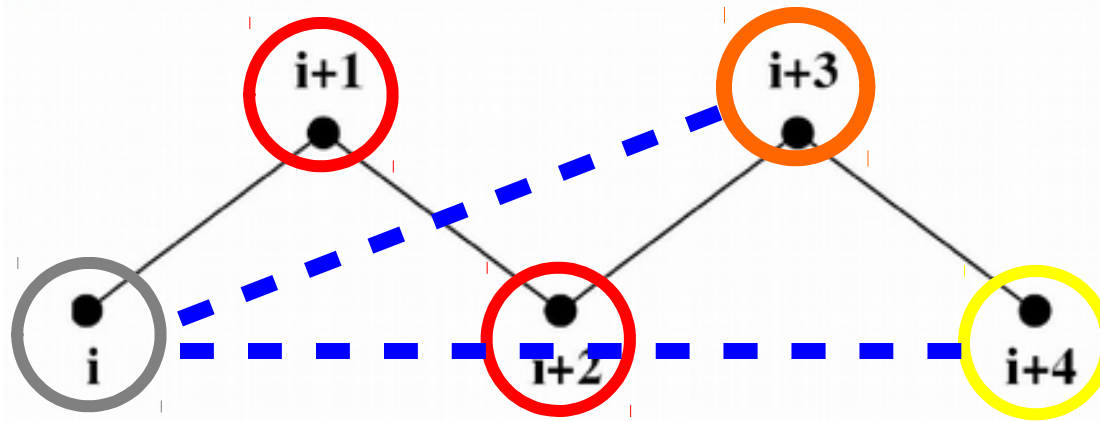
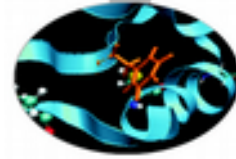


Figure 4.15: Atoms along an alkane chain.

Atoms covalently bound are defined as first neighbours second neighbours and so on....

LJ and electrostatic interactions **are not calculated among** first and second neighbours since they are considered in the stretching (first) or in the bending potential (second)

The standard non-bonding interactions are too strong for the third neighbours and are reduced (interactions 1-4; list 1-4)